

wi4243AP: Complex Analysis

week 7, Monday

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Outline

- 1 Section 6.3: Evaluating integrals
- 2 Section 6.4: Fourier transforms
 - Definitions and properties
 - Calculations
- 3 From the PDE course
- 4 Section 7.2: Laplace Transform
 - The transform
 - Inversion

One more example: $\int_0^{\infty} \frac{\ln x}{1+x^4} dx$

We take

$$f(z) = \frac{\log z}{1+z^4}$$

where $\log z$ is the branch of the logarithm defined by

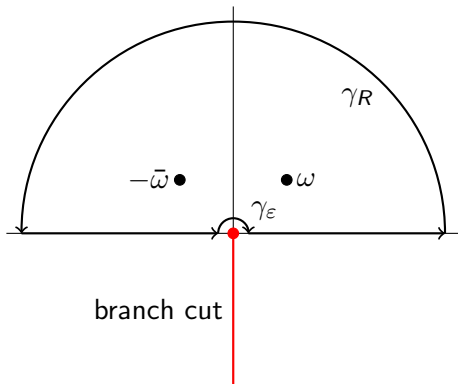
$$\log z = \ln |z| + i \arg z \quad \left(-\frac{1}{2}\pi < \arg z < \frac{3}{2}\pi\right)$$

so the branch cut is the negative imaginary axis.

Note:

- if $x > 0$ then $\log x = \ln x$;
- if $x < 0$ then $\log x = \ln |x| + i\pi$.

$\int_0^{\infty} \frac{\ln x}{1+x^4} dx$: the contour



The contour $\Gamma_{\epsilon, R}$ consists of

Interval $[\epsilon, R]$

Semicircle of radius R
(anticlockwise)

Interval $[-R, -\epsilon]$

Semicircle of radius ϵ (clockwise)

and the singularities

$\int_0^{\infty} \frac{\ln x}{1+x^4} dx$: singularities and residues

The singularities of f are the zeros of $z^4 + 1$, to wit ω , $-\omega$, $\bar{\omega}$ and $-\bar{\omega}$, where

$$\omega = e^{\frac{\pi}{4}i} = \frac{1}{2}\sqrt{2} + \frac{i}{2}\sqrt{2}.$$

of these ω and $-\bar{\omega}$ lie inside $\Gamma_{\varepsilon,R}$ (when $\varepsilon < 1 < R$ of course); the residues are

$$\operatorname{Res}(f, \omega) = \frac{\frac{1}{4}\pi i}{4i\omega} = \frac{\pi}{16}\bar{\omega} \quad \text{and} \quad \operatorname{Res}(f, -\bar{\omega}) = \frac{\frac{3}{4}\pi i}{4i\bar{\omega}} = \frac{3\pi}{16}\omega$$

so

$$\oint_{\Gamma_{\varepsilon,R}} f(z) dz = 2\pi i \cdot \frac{\pi}{16}(\bar{\omega} + 3\omega) = -\frac{\pi^2\sqrt{2}}{8} + \frac{\pi^2\sqrt{2}}{4}i$$

$\int_0^\infty \frac{\ln x}{1+x^4} dx$: Splitting the integral

The integral splits into

- $\int_\epsilon^R \frac{\ln x}{1+x^4} dx$
- $\int_{\gamma_R} f(z) dz$
- $\int_{-R}^{-\epsilon} \frac{\ln|x| + \pi i}{1+x^4} dx = \int_\epsilon^R \frac{\ln x}{1+x^4} dx + \pi i \int_\epsilon^R \frac{1}{1+x^4} dx$
- $\int_{\gamma_\epsilon} f(z) dz$

$\int_0^\infty \frac{\ln x}{1+x^4} dx$: Estimates

The big semicircle:

$$\left| \oint_{\gamma_R} f(z) dz \right| \leq \pi R \frac{\ln R + \pi}{R^4 - 1}$$

the right-hand side converges to 0 as $R \rightarrow \infty$

The small semicircle

$$\left| \oint_{\gamma_\varepsilon} f(z) dz \right| \leq \pi \varepsilon \frac{-\ln \varepsilon + \pi}{1 - \varepsilon^4}$$

the right-hand side converges to 0 as $\varepsilon \rightarrow 0$

$\int_0^\infty \frac{\ln x}{1+x^4} dx$: Putting it together

After we let $R \rightarrow \infty$ and $\varepsilon \rightarrow 0$ we get

$$2 \int_0^\infty \frac{\ln x}{1+x^4} dx + \pi i \int_0^\infty \frac{1}{1+x^4} dx = -\frac{\pi^2 \sqrt{2}}{8} + \frac{\pi^2 \sqrt{2}}{4} i$$

so that

$$\int_0^\infty \frac{\ln x}{1+x^4} dx = -\frac{\pi^2 \sqrt{2}}{16}$$

and

$$\int_0^\infty \frac{1}{1+x^4} dx = \frac{\pi \sqrt{2}}{4}$$

(which we knew already).

Definition

If $u : \mathbb{R} \rightarrow \mathbb{R}$ is such that $\int_{-\infty}^{\infty} |u(x)| dx < \infty$ then we put

$$\mathcal{F}\{u(x)\}(\omega) = \int_{-\infty}^{\infty} e^{i\omega x} u(x) dx \quad (\omega \text{ real})$$

the **Fourier transform** of u

(common abbreviation: $\mathcal{F}\{u(x)\}(\omega) = U(\omega)$).

Properties

linearity $\mathcal{F}\{\alpha u + \beta v\} = \alpha \mathcal{F}\{u\} + \beta \mathcal{F}\{v\}$

shifting $\mathcal{F}\{u(x - a)\} = e^{ia\omega} \mathcal{F}\{u(x)\}$

shifting $\mathcal{F}\{e^{-i\omega_0 x} u(x)\} = U(\omega - \omega_0)$

scaling $\mathcal{F}\{u(bx)\} = \frac{1}{b} U\left(\frac{1}{b}\omega\right)$

derivatives $\mathcal{F}\{u'(x)\} = -i\omega \mathcal{F}\{u(x)\}$

derivatives $\frac{d}{d\omega} \mathcal{F}\{u(x)\} = i\omega \mathcal{F}\{u(x)\}$

convolution $\mathcal{F}\{u * v\} = \mathcal{F}\{u\} \cdot \mathcal{F}\{v\}$

Convolution product: $(u * v)(x) = \int_{-\infty}^{\infty} u(\xi)v(x - \xi) d\xi$

Inversion and Parseval's identity

The Fourier transform is invertible:

$$u(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{-i\omega x} U(\omega) d\omega$$

It preserves norms:

$$\int_{-\infty}^{\infty} |u(x)|^2 dx = \frac{1}{2\pi} \int_{-\infty}^{\infty} |U(\omega)|^2 d\omega$$

Warning

Four possible conventions:

- $\int_{-\infty}^{\infty} e^{i\omega x} u(x) dx$ (the present text)
- $\frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{i\omega x} u(x) dx$ (norm-preserving)
- $\int_{-\infty}^{\infty} u(x) e^{-i\omega x} dx$ (looks like complex inner product)
- $\frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} u(x) e^{-i\omega x} dx$ (mathematician's favourite)

Just remember to check the convention of the book that you are reading.

\mathcal{F} as an excuse to calculate residues

We can calculate

$$\int_{-\infty}^{\infty} e^{i\omega x} f(x) dx$$

if f extends to a 'tame' complex function.

- a finite number of isolated singularities
- $\lim_{z \rightarrow \infty} f(z) = 0$
- no singularities on the real axis

A useful lemma

Lemma (Jordan)

If $\lim_{z \rightarrow \infty} f(z) = 0$ and if $\omega > 0$ then

$$\lim_{R \rightarrow \infty} \int_{C_R} e^{i\omega z} f(z) dz = 0$$

where C_R is the semi-circle $\{z : |z| = r, \operatorname{Im} z \geq 0\}$.

Write $M_R = \max\{|f(z)| : z \in C_R\}$; by assumption

$$\lim_{R \rightarrow \infty} M_R = 0$$

Jordan's Lemma: proof

Parametrize the semi-circle: $z = Re^{it}$ ($0 \leq t \leq \pi$) and rewrite the integral:

$$\int_0^\pi f(Re^{it}) e^{i\omega R \cos t} e^{-\omega R \sin t} Re^{it} dt$$

From this we get:

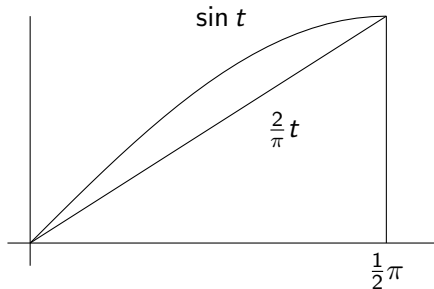
$$\begin{aligned} \left| \int_{C_R} f(z) dz \right| &\leq M_R \cdot R \cdot \int_0^\pi e^{-\omega R \sin t} dt \\ &\stackrel{*}{\leq} 2M_R \cdot R \cdot \int_0^{\frac{1}{2}\pi} e^{-\omega R \frac{2}{\pi} t} dt \\ &= 2M_R \cdot R \cdot \frac{\pi}{2\omega R} (1 - e^{-\omega R}) \end{aligned}$$

Jordan's Lemma: proof

Clearly:

$$\lim_{R \rightarrow \infty} \frac{M_R \pi}{\omega} (1 - e^{-\omega R}) = 0$$

And inequality * follows from $\sin t \geq \frac{2}{\pi} t$ ($0 \leq t \leq \frac{1}{2}\pi$)



Consequences

If $\omega > 0$ then

$$\mathcal{F}\{f(x)\} = 2\pi i [\text{sum of residues of } e^{i\omega z} f(z) \text{ above real axis}]$$

If $\omega < 0$ then we can take the semi-circle below the real axis to obtain

$$\mathcal{F}\{f(x)\} = -2\pi i [\text{sum of residues of } e^{i\omega z} f(z) \text{ below real axis}]$$

Example: $\mathcal{F}\left\{\frac{1}{x^4+1}\right\}$

Take $\omega > 0$ and rename the zeros of $z^4 + 1$ as α , $\bar{\alpha}$, $-\alpha$ and $-\bar{\alpha}$ (still: $\alpha = \frac{1}{2}\sqrt{2} + \frac{i}{2}\sqrt{2}$).

From earlier examples we know what we will get:

$$\text{Res}(e^{i\omega z} f(z), \alpha) = \frac{1}{4i} \bar{\alpha} e^{i\omega \alpha} \quad \text{and} \quad \text{Res}(e^{i\omega z} f(z), -\bar{\alpha}) = \frac{1}{4i} \alpha e^{-i\omega \bar{\alpha}}$$

After some 'hard' work:

$$\begin{aligned} \mathcal{F}\left\{\frac{1}{x^4+1}\right\} &= \frac{2\pi i}{4i} (\bar{\alpha} e^{i\omega \alpha} + \alpha e^{-i\omega \bar{\alpha}}) \\ &= \frac{\pi}{2} \sqrt{2} e^{-\frac{\omega}{2}\sqrt{2}} (\cos \frac{\omega}{2}\sqrt{2} + \sin \frac{\omega}{2}\sqrt{2}) \end{aligned}$$

One-dimensional Klein-Gordon equation

To find its Green's function solve

$$\frac{\partial^2 g}{\partial x^2} - \frac{1}{c^2} \left(\frac{\partial^2 g}{\partial t^2} + \alpha^2 g \right) = -\delta(x - \xi)\delta(t - \tau)$$

with $-\infty < x, \xi < \infty$ and $t, \tau > 0$ and where c is a positive constant.

Initial conditions: $g(x, 0|\xi, \tau) = g_t(x, 0|\xi, \tau) = 0$

Boundary conditions: $\lim_{|x| \rightarrow \infty} |g(x, t|\xi, \tau)| < \infty$.

Transform this

First Laplace ($t \rightarrow s$):

$$\frac{\partial^2 G}{\partial x^2} - \left(\frac{s^2 + a^2}{c^2} \right) G = -\delta(x - \xi) e^{-s\tau}$$

then Fourier ($x \rightarrow \omega$):

$$\left(-\omega^2 - \frac{s^2 + a^2}{c^2} \right) \mathcal{F}\{G\} = -e^{-s\tau} e^{i\omega\xi}$$

or

$$\mathcal{F}\{G\} = \frac{e^{-s\tau} e^{i\omega\xi}}{\omega^2 + \frac{1}{c^2}(s^2 + a^2)}$$

Inversion

So now:

$$\begin{aligned} G(x, s|\xi, \tau) &= \frac{1}{2\pi} e^{-s\tau} \int_{-\infty}^{\infty} e^{-i\omega x} \frac{e^{i\omega\xi}}{\omega^2 + \frac{1}{c^2}(s^2 + a^2)} d\omega \\ &= \frac{1}{2\pi} e^{-s\tau} \int_{-\infty}^{\infty} \frac{e^{-i(x-\xi)\omega}}{\omega^2 + \frac{1}{c^2}(s^2 + a^2)} d\omega \end{aligned}$$

Singularities and residues

The singularities are at $\omega^+ = \frac{i}{c}\sqrt{s^2 + a^2}$ and $\omega^- = -\frac{i}{c}\sqrt{s^2 + a^2}$,
with residues

$$\text{Res}(\cdot, \omega^+) = \frac{e^{\frac{x-\xi}{c}\sqrt{s^2+a^2}}}{2\frac{i}{c}\sqrt{s^2+a^2}}$$

and

$$\text{Res}(\cdot, \omega^-) = \frac{e^{-\frac{x-\xi}{c}\sqrt{s^2+a^2}}}{-2\frac{i}{c}\sqrt{s^2+a^2}}$$

The integrals

From the formula remember that we take the residues above the real axis if $-(x - \xi) > 0$ (i.e., $x < \xi$) and those below the real axis if $-(x - \xi) < 0$ (i.e., $x > \xi$).

So if $x < \xi$ we get

$$\begin{aligned} G(x, s|\xi, \tau) &= \frac{1}{2\pi} e^{-s\tau} \cdot 2\pi i \operatorname{Res}(\cdot, \omega^+) \\ &= \frac{ce^{-s\tau}}{2} \frac{\exp\left(\frac{x-\xi}{c} \sqrt{s^2 + a^2}\right)}{\sqrt{s^2 + a^2}} \end{aligned}$$

The integrals

And if $x > \xi$ we get

$$\begin{aligned} G(x, s|\xi, \tau) &= \frac{1}{2\pi} e^{-s\tau} \cdot -2\pi i \operatorname{Res}(\cdot, \omega^-) \\ &= \frac{ce^{-s\tau} \exp\left(-\frac{x-\xi}{c} \sqrt{s^2 + a^2}\right)}{2 \sqrt{s^2 + a^2}} \end{aligned}$$

Note: the minus sign in front of the $2\pi i$ is because we integrate **clockwise** around the singularity in the lower half plane.

A single formula

We can combine both answers into one:

$$G(x, s|\xi, \tau) = \frac{ce^{-s\tau} \exp\left(-\frac{|x-\xi|}{c} \sqrt{s^2 + a^2}\right)}{2 \sqrt{s^2 + a^2}}$$

Definition

The formula

For $f : [0, \infty) \rightarrow \mathbb{C}$ its Laplace transform is given by

$$F(s) = \mathcal{L}\{f(t)\} = \int_0^{\infty} e^{-st} f(t) dt$$

Here $s = \sigma + i\tau$ is complex.

If there are M and a such that $|f(t)| \leq Me^{at}$ for $t \geq 0$ then $\mathcal{L}\{f(t)\}$ exists whenever $\operatorname{Re} s > a$.

Properties

The map \mathcal{L} is linear and injective, i.e.,

- $\mathcal{L}\{\alpha f + \beta g\} = \alpha\mathcal{L}\{f\} + \beta\mathcal{L}\{g\}$ (where defined)
- if $\mathcal{L}\{f\} = \mathcal{L}\{g\}$ for σ large enough then $f = g$

Also, Leibniz' rule applies:

$$\frac{d}{ds} \int_0^{\infty} e^{-st} f(t) dt = \int_0^{\infty} \frac{\partial}{\partial s} e^{-st} f(t) dt = \int_0^{\infty} -te^{-st} f(t) dt$$

so, $F'(s) = \mathcal{L}\{-t \cdot f(t)\}$ and $F(s)$ is analytic.

Example: $\sin t$

We calculate $\mathcal{L}\{\sin t\}$:

$$\frac{1}{2i}(\mathcal{L}\{e^{it}\} - \mathcal{L}\{e^{-it}\}) = \frac{1}{2i} \left(\frac{1}{s-i} - \frac{1}{s+i} \right) = \frac{1}{s^2 + 1}$$

Note

- $\int_0^\infty e^{-st} \sin t \, dt$ exists whenever $\operatorname{Re} s > 0$
- the resulting function has an analytic extension to all of \mathbb{C} (except at i and $-i$).

Luckily this happens a lot, see the formulas on pp 327–9.

Bromwich integral

The formula

Assume $|f(t)| \leq Me^{\alpha t}$ for some α and let $a > \alpha$. Then

$$f(t) = \frac{1}{2\pi i} \lim_{R \rightarrow \infty} \int_{a-iR}^{a+iR} e^{st} F(s) ds$$

the integral is along the vertical straight line $\operatorname{Re} s = a$.

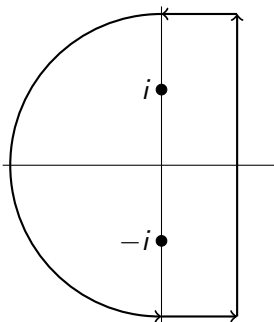
So, e.g.,

$$\sin t = \frac{1}{2\pi i} \lim_{R \rightarrow \infty} \int_{a-iR}^{a+iR} e^{st} \frac{1}{s^2 + 1} ds$$

Example: $\sin t$

Let's calculate the integral for $\sin t$.

We use the following contour Γ_R with two singularities inside



The segment $[a - iR, a + iR]$

The segment $[a + iR, iR]$

The semi-circle γ_R

The segment $[-iR, a - iR]$

The contour integral

Of course

$$\frac{1}{2\pi i} \oint_{\Gamma_R} \frac{e^{st}}{s^2 + 1} ds = \text{Res}(f, i) + \text{Res}(f, -i) = \frac{e^{it}}{2i} + \frac{e^{-it}}{-2i} = \sin t$$

So, all we need to do is show that

$$\lim_{R \rightarrow \infty} \frac{1}{2\pi i} \oint_{\Gamma_R} \frac{e^{st}}{s^2 + 1} ds = \lim_{R \rightarrow \infty} \int_{a-iR}^{a+iR} \frac{e^{st}}{s^2 + 1} ds$$

The other bits

$[a + iR, iR]$ length: a , upper bound for modulus: $e^{at}/(R^2 - 1)$

$[-iR, a - iR]$ length: a , upper bound for modulus: $e^{at}/(R^2 - 1)$

γ_R length: πR , upper bound for modulus: $1/(R^2 - 1)$

In all three cases: limit of integral is zero.

Done.

What to do?

From the book: 6.4, 7.2, transforms resulting from PDE
Suitable problems: 6.23 - 6.28; 7.16 - 7.25.
Recommended problems: 6.23, 6.27, 7.22, 7.23