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Abstract - Let E be a separable real Banach space and let $Q \in \mathcal{L}(E^*, E)$ be positive and symmetric. Let $\mathbf{S} = \{S(t)\}_{t \ge 0}$ be a C_0 -semigroup on E. We study the relations between the reproducing kernel Hilbert spaces associated with the operators $Q_t := \int_0^t S(s)QS^*(s) \, ds$. Under the assumption that these operators are the covariances of centered Gaussian measures μ_t on E, we also study equivalence $\mu_t \sim \mu_s$ for different values of s and t, and we calculate their Radon-Nikodym derivatives.

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0. Introduction

In this paper we investigate the reproducing kernel Hilbert spaces and Gaussian measures associated with a non-symmetric Ornstein-Uhlenbeck semigroup on a separable real Banach space E. This study is usually carried out in a Hilbert space setting, and one of the motivations of this paper was to see to what extent the theory can be extended to the Banach space setting.

The main difference between the Banach space- and the Hilbert space situation is that the covariance operator of a Gaussian measure on a Banach space E is a positive symmetric operator Q (the precise definitions are given in Section 1) from the dual E^* into E, rather than an operator on E. Thus, in contrast to the Hilbert space situation, it is no longer possible to represent the reproducing kernel Hilbert space H associated with Q as $H = \text{Im } Q^{1/2}$. When working in a Banach space setting, any reference to the operator $Q^{1/2}$ has therefore to be avoided. This turns out to be, at least for the

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questions considered in this paper, more of an advantage than a disadvantage, as we believe that the resulting proofs have gained some transparency.

Another difference from the Hilbert space situation is that no necessary and sufficient conditions on a positive symmetric $Q \in \mathcal{L}(E^*, E)$ seem to be known in order that Q be the covariance operator of a Gaussian measure on E. As we will show, for several well-known results on non-symmetric Ornstein-Uhlenbeck semigroups it is not relevant whether or not the positive self-adjoint operators that one is led to, are covariances or not. In the remaining results we get around this difficulty by simply imposing that Q be the covariance of a Gaussian measure; this replaces the usual assumption in the Hilbert space setting that Q should be trace class.

Let us now describe in more detail the contents of this paper. Let E be a real Banach space, let $Q \in \mathcal{L}(E^*, E)$ be positive and symmetric, and let $\mathbf{S} = \{S(t)\}_{t \ge 0}$ be a C_0 -semigroup on E. The operators

$$Q_t := \int_0^t S(s)QS^*(s) \, ds$$

are well-defined in the strong sense (cf. Proposition 1.2 below), and positive and symmetric. In case E is a Hilbert space and Q_t is also trace class, Q_t can be identified as the covariance of the distribution μ_t of the E-valued Gaussian random variable

$$X(t) = \int_0^t S(t-s) \, dW_Q(s),$$

where W_Q denotes a cylindrical Q-Wiener process and the integral is an Itô type stochastic integral. The importance of this resides in the fact that the (strong Markov) process $(X(t))_{t\geq 0}$ is the unique weak solution of the Langevin equation

$$dX(t) = AX(t) dt + dW_Q(t), \qquad t \ge 0,$$

$$X(0) = 0 \quad \text{almost surely,}$$

where A is the infinitesimal generator of **S**. Without the trace class assumption similar results hold; this time W_Q has to be interpreted as a cylindrical Q-Wiener process. For a comprehensive treatment of these concepts we refer to the book [DZ3].

In Section 1 we undertake a detailed study of the reproducing kernel Hilbert spaces (RKHS's) H_t associated with the operators Q_t . We do not assume that Q_t be the covariance operators of Gaussian measures on E. We prove that

$$S(s)H_{t_0} \subset H_{t_0+s}$$

for all s > 0 and $t_0 > 0$, and that $H_{t_0} = H_{t_0+s}$ (as subsets of E) if and only if $S(t_0)$, regarded as an operator from H_s into H_{t_0+s} , is a strict contraction. We also show that

$$S(s)H_{t_0} \subset H_{t_0}$$

for given s > 0 and $t_0 > 0$ implies

$$H_t = H_{\max\{s, t_0\}} \quad \text{for all} \quad t \in [\max\{s, t_0\}, \infty)$$

and that this result is the best possible.

In Sections 2 through 5 we assume that E is separable and that each of the Q_t is the covariance operator of a centered Gaussian measures μ_t on E. After some preliminary observations in Section 2, we study equivalence of measures $\mu_t \sim \mu_s$ under various conditions in Section 3. For instance, it is shown that

$$\mu_t \sim \mu_{t_0}$$
 for all $t \in [t_0, \infty)$

whenever there exist $s \in (0, \infty)$ and $t_1 \in (t_0, \infty)$ such that $S(s)Q = QS^*(s)$ and $\mu_{t_1} \sim \mu_{t_0}$.

In Section 4 we derive an explicit formula for the Radon-Nikodym derivative $d\mu_{t_0}/d\mu_{t_1}$ whenever these measures are equivalent. The approach depends on second quantization, existence of linear μ -measurable extensions, and a classical theorem of Shale concerning absolute continuity of image measures, and may be of some interest in its own right.

In Section 5 we proceed to show that for t_1 fixed the Radon-Nikodym derivative $d\mu_{t_0}/d\mu_{t_1}$ depends continuously upon t_0 .

In Section 6 we return to the cylindrical case and study the RKHS H_{∞} associated with the strong limit $Q_{\infty} = \lim_{t\to\infty} Q_t$ whenever this limit exists. Assuming that Q_{∞} is the covariance of an (invariant) centered Gaussian measure μ_{∞} on E, we discuss versions for μ_{∞} of some of the results obtained in the previous sections. For Hilbert spaces E, the main results of this section have been obtained recently by Chojnowska-Michelak and Goldys [CG1-3], [Go] and Fuhrman [Fu]. In particular this is true for the expression of the Radon-Nikodym derivative $d\mu_t/d\mu_{\infty}$, which was established by [Fu] in the null controllable case, and was extended to the more general situation considered here by [CG3]. We point out, however, that the approach taken in these references in very different for ours. To the best of our knowledge the principal results of Sections 1 through 5 are new even in the Hilbert space setting. Some of these (Theorems 1.4, 1.7, 3.2, and 4.1) extend in a natural corresponding results about invariant measures to finite t, but others have no analogue for invariant measures (Theorems 1.9 and 3.5) or its analogue seems to be new as well (Theorem 5.5).

In the final Section 7 we discuss some extensions of our results to the class of so-called (cylindrical) Gaussian Mehler semigroups recently introduced by Bogachev, Röckner, and Schmuland [BRS].

1. The reproducing kernel Hilbert spaces H_t

Throughout this section, E is a fixed arbitrary real Banach space. A bounded linear operator $Q \in \mathcal{L}(E^*, E)$ is called *positive* if $\langle Qx^*, x^* \rangle \ge 0$ for all $x^* \in E^*$ and symmetric

if $\langle Qx^*, y^* \rangle = \langle Qy^*, x^* \rangle$ for all $x^*, y^* \in E^*$. If Q is positive and symmetric, then on $\operatorname{Im} Q = \{Qx^* : x^* \in E^*\}$ we may define an inner product $[\cdot, \cdot]$ by the formula $[Qx^*, Qy^*] := \langle Qx^*, y^* \rangle$. The completion H of $\operatorname{Im} Q$ with respect to this inner product is a Hilbert space, and the inclusion $i : \operatorname{Im} Q \subset E$ extends to a continuous injection $i : H \to E$. Moreover, if we regard Q as an operator from E^* to H we have the identity $i^* = Q$. We will refer to H as the *reproducing kernel Hilbert space* (RKHS) associated with Q. If E is separable, then H is separable as well. If E is a *Hilbert* space and Q is a positive and symmetric operator on E (identifying E and its dual), then $H = \operatorname{Im} Q^{\frac{1}{2}}$ with identical norms. For more information we refer to [VTC, Chapter III], where the simple proofs can be found.

We recall with the following result, which is proved along the lines of [DZ2, Proposition B.1].

Proposition 1.1. Let $Q, \tilde{Q} \in \mathcal{L}(E^*, E)$ be two positive symmetric operators. Then for the associated RKHS's we have $H \subset \tilde{H}$ (as subsets of E) if and only if there exists a constant K > 0 such that

$$\langle Qx^*, x^* \rangle \leqslant K \langle \tilde{Q}x^*, x^* \rangle, \qquad \forall x^* \in E^*.$$

In this situation, the operator $V : \operatorname{Im} \tilde{Q} \to \operatorname{Im} Q$ defined by $V\tilde{Q}x^* := Qx^*$ extends to a bounded operator from \tilde{H} into H; we will sometimes use the suggestive notation $V = Q\tilde{Q}^{-1}$. If $H = \tilde{H}$ (as subsets of E), V is a (Banach space) isomorphism of \tilde{H} onto H, the inverse being given by $V^{-1}Qx^* = \tilde{Q}x^*$.

Suppose $Q \in \mathcal{L}(E^*, E)$ is positive and symmetric, and let $\mathbf{S} = \{S(t)\}_{t \ge 0}$ be a C_0 -semigroup on E. Our terminology concerning C_0 -semigroups is standard; we refer to [Pa] for more details. For each t > 0 the operator Q_t defined by

$$Q_t x^* := \int_0^t S(s) Q S^*(s) x^* \, ds, \qquad x^* \in E^*,$$

is positive and symmetric. Note that this integral exists as a Bochner integral in E; strong measurability of the integrand follows from:

Proposition 1.2. For all $x^* \in E^*$, the function $s \mapsto S(s)QS^*(s)x^*$ is strongly measurable.

Proof: As a map from E^* into H, the operator Q is the adjoint of the inclusion map $i: H \subset E$, and as such Q is weak*-to-weakly continuous. Hence the weak*-continuity of $S(\cdot)x^*$ implies weak continuity of $QS^*(\cdot)x^*$.

Step 1 - First we assume that E is separable. Then H is separable and we may choose a countable orthonormal basis $(h_n) \subset H$. Fix $y^* \in E^*$. Expanding $QS^*(s)x^*$ and $QS^*(s)y^*$ in terms of (h_n) we have

$$\langle S(s)QS^*(s)x^*, y^* \rangle = [QS^*(s)x^*, QS^*(s)y^*] = \sum_{n=1}^{\infty} [QS^*(s)x^*, h_n][QS^*(s)y^*, h_n],$$

so $\langle S(\cdot)QS^*(\cdot)x^*, y^* \rangle$ appears as a countable sum of continuous functions. This proves that $S(\cdot)QS^*(\cdot)x^*$ is weakly measurable. Since it is also separably valued by the separability of E, strong measurability now follows by an appeal to Pettis's measurability theorem [DU, Chapter II].

Step 2 - Now let E be arbitrary. Let H_0 be the closed linear span in H of the set $\{QS^*(t)x^*: t \ge 0\}$. Since Q, as a map from E^* to H, is weak*-to-weakly continuous and $t \mapsto S^*(t)x^*$ is weak*-continuous, H_0 is weakly separable and therefore separable. Denoting by E_0 the smallest closed **S**-invariant subspace in E containing H_0 , it follows that E_0 is separable in E. Let $i_0 : H_0 \subset E_0$ and $j_0 : E_0 \subset E$ denote the inclusion maps. Now define $Q_0 \in \mathcal{L}(E_0^*, E_0)$ by

$$Q_0(j_0^*y^*) := (i_0 \circ P_0 \circ Q)y^*, \qquad y^* \in E^*,$$

where P_0 is the orthogonal projection of H onto H_0 . We check that Q_0 is well-defined. If $j_0^* y^* = 0$, then y^* annihilates E_0 and therefore, for all $t \ge 0$,

$$[QS^*(t)x^*, Qy^*] = \langle QS^*(t)x^*, y^* \rangle = 0.$$

This means that $Qy^* \perp H_0$, so $P_0Qy^* = 0$ and hence $Q_0(j_0^*y^*) = 0$. Next we check that Q_0 is positive and symmetric. For all $y^* \in E^*$ and $z^* \in E^*$ we have

$$\langle Q_0 j_0^* y^*, j_0^* z^* \rangle = \langle i P_0 Q y^*, z^* \rangle = [P_0 Q y^*, Q z^*] = [P_0 Q y^*, P_0 Q z^*],$$

which is symmetric in y^* and z^* and non-negative if $y^* = z^*$.

Let \mathbf{S}_0 denote the restriction of \mathbf{S} to the invariant subspace E_0 . The lemma follows from the corresponding result for the separable space E_0 once we have realized that for all $s \ge 0$,

$$S(s)QS^*(s)x^* = S_0(s)i_0P_0QS^*(s)x^* = S_0(s)Q_0(j_0^*S^*(s)x^*) = S_0(s)Q_0S_0^*(s)(j_0^*x^*).$$

We will frequently use the following algebraic relation between the operators Q_t , which is immediate from their definition: for all t, s > 0 we have

$$Q_{t+s} = Q_s + S(s)Q_t S^*(s).$$

The RKHS associated with Q_t will be denoted by H_t ; the inclusion map $H_t \subset E$ is denoted by i_t . As in the case of a Hilbert space E, H_t can be interpreted as the space of reachable states of a certain linear control problem in E; this point of view will be elaborated elsewhere.

The present section is devoted to a systematic study of the relation between the spaces H_t for different values of t. The first observation is a direct consequence of Proposition 1.1:

Proposition 1.3. If $0 < t_0 \leq t_1$, then $H_{t_0} \subset H_{t_1}$.

From the identity $S(s)Q_{t_0}S^*(s) = Q_{t_0+s} - Q_s$ combined with Proposition 1.3 we see that S(s) maps the linear subspace $\operatorname{Im}(Q_{t_0}S^*(s))$ of H_{t_0} into H_{t_0+s} . The next result shows that we actually have $S(s)H_{t_0} \subset H_{t_0+s}$:

Theorem 1.4. For all s > 0 and $t_0 > 0$ we have $S(s)H_{t_0} \subset H_{t_0+s}$. Moreover, $||S(s)||_{\mathcal{L}(H_{t_0},H_{t_0+s})} \leq 1$.

Proof: For all $x^* \in E^*$ we have

$$\begin{aligned} \|Q_{t_0}S^*(s)x^*\|_{H_{t_0}}^2 &= \langle Q_{t_0}S^*(s)x^*, S^*(s)x^* \rangle \\ &= \langle Q_{t_0+s}x^*, x^* \rangle - \langle Q_sx^*, x^* \rangle \\ &\leqslant \langle Q_{t_0+s}x^*, x^* \rangle = \|Q_{t_0+s}x^*\|_{H_{t_0+s}}^2. \end{aligned}$$
(1.1)

Hence,

$$|\langle Q_{t_0}S^*(s)x^*, y^*\rangle| = |[Q_{t_0}S^*(s)x^*, Q_{t_0}y^*]_{H_{t_0}}| \leq ||Q_{t_0+s}x^*||_{H_{t_0+s}} ||Q_{t_0}y^*||_{H_{t_0}}.$$
 (1.2)

Define a linear functional ψ_{s,y^*} : Im $Q_{t_0+s} \to \mathbb{R}$ by

$$\psi_{s,y^*}(Q_{t_0+s}x^*) := \langle Q_{t_0}S^*(s)x^*, y^* \rangle.$$

If $Q_{t_0+s}x^* = 0$, then $Q_{t_0}S^*(s)x^* = 0$ by (1.1), so ψ_{s,y^*} is well-defined. By (1.2), ψ_{s,y^*} extends to a bounded linear functional on H_{t_0+s} of norm $\leq ||Q_{t_0}y^*||_{H_t}$. Identifying ψ_{s,y^*} with an element of H_{t_0+s} , for all $x^* \in E^*$ we have

$$\langle \psi_{s,y^*}, x^* \rangle = [Q_{t_0+s}x^*, \psi_{s,y^*}]_{H_{t_0+s}} = \langle Q_{t_0}S^*(s)x^*, y^* \rangle = \langle S(s)Q_{t_0}y^*, x^* \rangle.$$

Hence, $S(s)Q_{t_0}y^* = \psi_{s,y^*} \in H_{t_0+s}$ and $||S(s)Q_{t_0}y^*||_{H_{t_0+s}} \leq ||Q_{t_0}y^*||_{H_{t_0}}$.

Whenever it is convenient, the restriction of S(s) as an operator in $\mathcal{L}(H_t, H_{t+s})$ will be denoted by $S_{t\to t+s}(s)$ and its adjoint $(S_{t\to t+s}(s))^* \in \mathcal{L}(H_{t+s}, H_t)$ by $S_{t\to t+s}^*(s)$.

Corollary 1.5. For all $0 < t_0 < t_1$ the inclusion mapping $H_{t_0} \subset H_{t_1}$ is contractive. *Proof:* For all $x^* \in E^*$ we have

$$\begin{aligned} \|Q_{t_0}x^*\|_{H_{t_1}}^2 &= [Q_{t_0}x^*, Q_{t_1}x^* - S(t_0)Q_{t_1-t_0}S^*(t_0)x^*]_{H_{t_1}} \\ &= \langle Q_{t_0}x^*, x^* \rangle - [Q_{t_0}x^*, S(t_0)Q_{t_1-t_0}S^*(t_0)x^*]_{H_{t_1}} \\ &= \|Q_{t_0}x^*\|_{H_{t_0}}^2 - [Q_{t_0}x^*, S(t_0)Q_{t_1-t_0}S^*(t_0)x^*]_{H_{t_1}}. \end{aligned}$$

But $S_{t_1-t_0\to t_1}^*(t_0)Q_{t_1} = Q_{t_1-t_0}S^*(t_0)$. Hence,

$$\begin{split} &[Q_{t_0}x^*, S(t_0)Q_{t_1-t_0}S^*(t_0)x^*]_{H_{t_1}} \\ &= [Q_{t_1}x^*, S(t_0)Q_{t_1-t_0}S^*(t_0)x^*]_{H_{t_1}} - \|S(t_0)Q_{t_1-t_0}S^*(t_0)x^*\|_{H_{t_1}}^2 \\ &= \langle Q_{t_1-t_0}S^*(t_0)x^*, S^*(t_0)x^* \rangle - \|S(t_0)Q_{t_1-t_0}S^*(t_0)x^*\|_{H_{t_1}}^2 \\ &= \|Q_{t_1-t_0}S^*(t_0)x^*\|_{H_{t_1-t_0}}^2 - \|S_{t_1-t_0\to t_0}(t_0)Q_{t_1-t_0}S^*(t_0)x^*\|_{H_{t_1}}^2 \\ &\ge 0; \end{split}$$

for the inequality we used that $\|S_{t_1-t_0\to t_1}(t_0)\|_{\mathcal{L}(H_{t_1-t_0},H_{t_1})} \leq 1$. We conclude that $\|Q_{t_0}x^*\|_{H_{t_1}} \leq \|Q_{t_0}x^*\|_{H_{t_0}}$ for all $x^* \in E^*$, and the corollary follows.

Next we characterize equality of H_{t_0} and H_{t_0+s} in terms of the restriction $S(t_0) \in \mathcal{L}(H_{t_0}, H_{t_0+s})$. For later reference, we first isolate a simple lemma.

Lemma 1.6. Let $t_1 \ge t_0 > 0$, s > 0, and assume that S(s) maps H_{t_0} into H_{t_1} . Then $S(s) \in \mathcal{L}(H_{t_0}, H_{t_1})$, and for all $x^* \in E^*$ we have

$$\|Q_{t_0}S^*(s)x^*\|_{H_{t_0}} \leq \|S(s)\|_{\mathcal{L}(H_{t_0},H_{t_1})} \cdot \|Q_{t_1}x^*\|_{H_{t_1}}$$

Proof: By the closed graph theorem, S(s) is bounded as an operator from H_{t_0} into H_{t_1} . For all $x^* \in E^*$ we have

$$\begin{split} \|Q_{t_0}S^*(s)x^*\|_{H_{t_0}} &= \sup\{[Q_{t_0}S^*(s)x^*, Q_{t_0}y^*]_{H_{t_0}}: \ y^* \in E^*, \|Q_{t_0}y^*\|_{H_{t_0}} \leqslant 1\} \\ &= \sup\{\langle S(s)Q_{t_0}y^*, x^*\rangle: \ y^* \in E^*, \|Q_{t_0}y^*\|_{H_{t_0}} \leqslant 1\} \\ &= \sup\{[S(s)Q_{t_0}y^*, Q_{t_1}x^*]_{H_{t_1}}: \ y^* \in E^*, \|Q_{t_0}y^*\|_{H_{t_0}} \leqslant 1\} \\ &\leqslant \|S(s)\|_{\mathcal{L}(H_{t_0}, H_{t_1})} \cdot \|Q_{t_1}x^*\|_{H_{t_1}}. \end{split}$$

Theorem 1.7. Let $t_0 > 0$ and h > 0 be fixed. Then $H_{t_0} = H_{t_0+h}$ (as subsets of E) if and only if $||S(t_0)||_{\mathcal{L}(H_h, H_{t_0+h})} < 1$.

Proof: We have already seen that $H_{t_0} \subset H_{t_0+h}$. It remains to prove that $H_{t_0+h} \subset H_{t_0}$ if and only if $||S(t_0)||_{\mathcal{L}(H_h, H_{t_0+h})} < 1$.

First assume $||S(t_0)||_{\mathcal{L}(H_h, H_{t_0+h})} < 1$. For all $x^* \in E^*$ we have

$$\begin{aligned} \|Q_{t_0}x^*\|_{H_{t_0}}^2 &= \langle Q_{t_0}x^*, x^* \rangle \\ &= \langle Q_{t_0+h}x^*, x^* \rangle - \langle S(t_0)Q_hS^*(t_0)x^*, x^* \rangle \\ &= \|Q_{t_0+h}x^*\|_{H_{t_0+h}}^2 - \|Q_hS^*(t_0)x^*\|_{H_h}^2. \end{aligned}$$

But by Lemma 1.6,

$$\|Q_h S^*(t_0) x^*\|_{H_h} \leq \|S(t_0)\|_{\mathcal{L}(H_h, H_{t_0+h})} \cdot \|Q_{t_0+h} x^*\|_{H_{t_0+h}}.$$

Hence,

$$\langle Q_{t_0} x^*, x^* \rangle = \|Q_{t_0} x^*\|_{H_{t_0}}^2 \ge \left(1 - \|S(t_0)\|_{\mathcal{L}(H_h, H_{t_0+h})}^2\right) \|Q_{t_0+h} x^*\|_{H_{t_0+h}}^2$$
$$= \left(1 - \|S(t_0)\|_{\mathcal{L}(H_h, H_{t_0+h})}^2\right) \langle Q_{t_0+h} x^*, x^* \rangle.$$

By Proposition 1.1 this gives the inclusion $H_{t_0+h} \subset H_{t_0}$.

Conversely, assume that $H_{t_0+h} \subset H_{t_0}$. Then there exists K > 1 such that

 $\langle Q_{t_0+h}x^*, x^* \rangle \leqslant K \langle Q_{t_0}x^*, x^* \rangle = K \langle Q_{t_0+h}x^*, x^* \rangle - K \langle S(t_0)Q_hS^*(t_0)x^*, x^* \rangle$ for all $x^* \in E^*$, or equivalently,

 $\begin{aligned} \|Q_h S^*(t_0) x^*\|_{H_h}^2 &\leq (1 - K^{-1}) \|Q_{t_0 + h} x^*\|_{H_{t_0 + h}}^2 \\ \text{for all } x^* \in E^*. \text{ Hence for all } x^*, y^* \in E^*, \\ \|[S(t_0)Q_h y^*, Q_{t_0 + h} x^*]_{H_{t_0 + h}}\| &= \|[Q_h y^*, Q_h S^*(t_0) x^*]_{H_h}\| \\ &\leq \|Q_h y^*\|_{H_h} \|Q_h S^*(t_0) x^*\|_{H_h} \\ &\leq \sqrt{1 - K^{-1}} \|Q_h y^*\|_{H_h} \|Q_{t_0 + h} x^*\|_{H_{t_0 + h}}. \end{aligned}$

This shows that $||S(t_0)||_{\mathcal{L}(H_h, H_{t_0+h})} \leq \sqrt{1 - K^{-1}} < 1.$

Throughout the rest of this paper, the notation $H_{t_1} = H_{t_0}$ means equality of H_{t_1} and H_{t_0} as subsets of E; as Hilbert spaces, H_{t_1} and H_{t_0} will usually carry different inner products.

Corollary 1.8. If $0 < t_0 < t_1$ are such that $H_{t_1} = H_{t_0}$, then $H_t = H_{t_0}$ for all $t \in [t_0, \infty)$.

Proof: It is clear that $H_{t_0} = H_t = H_{t_1}$ for all $t \in [t_0, t_1]$. Furthermore, Theorem 1.7 implies that $H_{t_0+\delta} = H_{t_1+\delta}$ for all $\delta \ge 0$. These two observations clearly lead to the desired result.

The following theorem relates equality of different spaces H_t to their invariance under S:

Theorem 1.9. If $S(s)H_{t_0} \subset H_{t_0}$ for some s > 0, then $H_t = H_{\max\{s,t_0\}}$ for all $t \in [\max\{s,t_0\},\infty)$.

Proof: In view of the Proposition 1.3 we only need to prove the inclusion $H_t \subset H_{\max\{s,t_0\}}$ for $t \in [\max\{s,t_0\},\infty)$.

Step 1 - In this step we prove the following: If $\sigma \in (0, t_0]$ and $t_1 \in (t_0, 2t_0]$ are such that $S(\sigma)$ maps $H_{t_1-t_0}$ into H_{t_0} , then $H_{t_1} \subset H_{2t_0-\sigma}$. By Lemma 1.5, using that $0 < t_1 - t_0 \leq t_0$, for all $x^* \in E^*$ we have

$$\|Q_{t_1-t_0}S^*(\sigma)x^*\|_{H_{t_1-t_0}} \leq \|S(\sigma)\|_{\mathcal{L}(H_{t_1-t_0},H_{t_0})}\|Q_{t_0}x^*\|_{H_{t_0}}.$$

It follows that

$$\begin{aligned} \langle Q_{t_1}x^*, x^* \rangle &= \langle Q_{t_0}x^*, x^* \rangle + \langle Q_{t_1-t_0}S^*(t_0)x^*, S^*(t_0)x^* \rangle \\ &= \langle Q_{t_0}x^*, x^* \rangle + \|Q_{t_1-t_0}S^*(\sigma)S^*(t_0-\sigma)x^*\|^2_{H_{t_1-t_0}} \\ &\leqslant \langle Q_{t_0}x^*, x^* \rangle + \|S(\sigma)\|^2_{\mathcal{L}(H_{t_1-t_0}, H_{t_0})} \|Q_{t_0}S^*(t_0-\sigma)x^*\|^2_{H_{t_0}}. \end{aligned}$$

Now

$$\begin{aligned} \|Q_{t_0}S^*(t_0 - \sigma)x^*\|_{H_{t_0}}^2 &= \langle Q_{t_0}S^*(t_0 - \sigma)x^*, S^*(t_0 - \sigma)x^* \rangle \\ &= \langle Q_{2t_0 - \sigma}x^*, x^* \rangle - \langle Q_{t_0 - \sigma}x^*, x^* \rangle \\ &\leqslant \langle Q_{2t_0 - \sigma}x^*, x^* \rangle. \end{aligned}$$

Putting these estimates together, we obtain

$$\langle Q_{t_1}x^*, x^* \rangle \leqslant \langle Q_{t_0}x^*, x^* \rangle + \|S(\sigma)\|_{\mathcal{L}(H_{t_1-t_0}, H_{t_0})}^2 \langle Q_{2t_0-\sigma}x^*, x^* \rangle \leqslant \left(1 + \|S(\sigma)\|_{\mathcal{L}(H_{t_1-t_0}, H_{t_0})}^2\right) \langle Q_{2t_0-\sigma}x^*, x^* \rangle$$

By Proposition 1.1 this implies the inclusion $H_{t_1} \subset H_{2t_0-\sigma}$.

Step 2 - In this step we prove: If $s \in (0, t_0]$ is such that $S(s)H_{t_0} \subset H_{t_0}$, then for all $t_1 \in [t_0 + s, 2t_0]$ we have $H_{t_1} \subset H_{t_1-s}$. Indeed, by Theorem 1.4 we know that $S(2t_0 - t_1)$ maps $H_{t_1-t_0}$ into H_{t_0} . Therefore also $S(2t_0 - t_1 + s)H_{t_1-t_0} \subset H_{t_0}$, and from Step 1 we obtain $H_{t_1} \subset H_{2t_0-(2t_0-t_1+s)} = H_{t_1-s}$. Step 3 - In this step we prove the theorem for the case $s \in (0, t_0]$. First assume $t \in [t_0, 2t_0]$. Write $t = t_0 + ks + \varepsilon$, where k is a nonnegative integer and $\varepsilon \in [0, s)$. If k = 0, then by Proposition 1.3 and Step 2,

$$H_t \subset H_{t_0+s} \subset H_{t_0}.$$

If $k \ge 1$, then we apply Step 2 k times to see that

$$H_t \subset H_{t-s} \subset H_{t-2s} \subset \dots \subset H_{t-ks} = H_{t_0+\varepsilon},$$

and therefore by the previous case, $H_t \subset H_{t_0}$.

Step 4 - In this step we prove the theorem for the case $s \ge t_0$.

First observe that by dualizing the identity $i_{t_0}S(s)|_{H_{t_0}} = S(s)i_{t_0}$, where $i_{t_0} : H_{t_0} \subset E$ is the inclusion map, we obtain $(S(s)|_{H_{t_0}})^*Q_{t_0} = Q_{t_0}S^*(s)$. Fix $t_1 \in (s, s+t_0]$ arbitrary. For all $x^* \in E^*$ we have

$$\begin{split} \langle Q_{t_1}x^*, x^* \rangle &= \langle Q_s x^*, x^* \rangle + \langle Q_{t_1-s}S^*(s)x^*, S^*(s)x^* \rangle \\ &\leq \langle Q_s x^*, x^* \rangle + \langle Q_{t_0}S^*(s)x^*, S^*(s)x^* \rangle \\ &= \langle Q_s x^*, x^* \rangle + \|Q_{t_0}S^*(s)x^*\|_{H_{t_0}}^2 \\ &= \langle Q_s x^*, x^* \rangle + \|(S(s)|_{H_{t_0}})^*Q_{t_0}x^*\|_{H_{t_0}}^2 \\ &\leq \langle Q_s x^*, x^* \rangle + \|S(s)\|_{\mathcal{L}(H_{t_0})}^2 \langle Q_{t_0}x^*, x^* \rangle \\ &\leq \left(1 + \|S(s)\|_{\mathcal{L}(H_{t_0})}^2\right) \langle Q_s x^*, x^* \rangle. \end{split}$$

Hence, $H_{t_1} \subset H_s$. But then for any $\tau \in (0, t_1 - s]$, by Theorem 1.4 and Proposition 1.3 we have $S(\tau)H_s \subset H_{s+\tau} \subset H_{t_1} \subset H_s$. Since $0 < \tau \leq s$, Step 3 now shows that $H_t = H_s$ for all $t \geq s$.

Notice that the case $s = t_0$ already follows from Step 1. In Example 1.14 below we show that the bound max $\{s, t_0\}$ is the best possible.

Next we study the situation where H, the RKHS associated with Q, is **S**-invariant. Then by the closed graph theorem, for each t > 0 the restriction $S^H(t) := S(t)|_H$ is a bounded operator on H, and it is easy to see that the function $s \mapsto ||S^H(s)||_{\mathcal{L}(H)}$ is Borel.

Lemma 1.10. Suppose $S(t)H \subset H$ for all $t \ge 0$. If there exists T > 0 such that

$$\int_0^T \|S^H(s)\|_{\mathcal{L}(H)}^2 \, ds < \infty,$$

then $H_t \subset H$ for all t > 0.

Proof: By the semigroup property, for all t > 0 we have

$$\int_0^t \|S^H(s)\|_{\mathcal{L}(H)}^2 \, ds < \infty.$$

Then,

$$\begin{split} \langle Q_t x^*, x^* \rangle &= \int_0^t \langle QS^*(s)x^*, S^*(s)x^* \rangle \, ds \\ &= \int_0^t \|QS^*(s)x^*\|_H^2 \, ds \\ &= \int_0^t \|(S^H(s))^*Qx^*\|_H^2 \, ds \\ &\leqslant \langle Qx^*, x^* \rangle \int_0^t \|S^H(s)\|_{\mathcal{L}(H)}^2 \, ds, \end{split}$$

where we used that $iS^H(s) = S(s)i$ and $i^* = Q$ (recall that $i : H \subset E$ is the inclusion map) imply $(S^H(s))^*Q = QS^*(s)$. From Proposition 1.1 it follows that $H_t \subset H$.

Theorem 1.11. Suppose $S(t)H \subset H$ for all $t \ge 0$ and assume there exists T > 0 such that

$$\int_0^T \|S^H(s)\|_{\mathcal{L}(H)}^2 \, ds < \infty.$$

Then for each t > 0,

$$Q_t^H(Qx^*) := Q_t x^*, \qquad x^* \in E^*,$$

defines a bounded self-adjoint operator Q_t^H on H. Denoting the RKHS associated with the operator Q_t^H by \mathcal{H}_t , we have $H_t = \mathcal{H}_t$ with identical norms.

Proof: For all $x^* \in E^*$ and $y^* \in E^*$ we have

$$[Q_t^H(Qx^*), Qy^*]_H = \int_0^t [S^H(s)QS^*(s)x^*, Qy^*]_H \, ds$$
$$= \int_0^t [S^H(s)(S^H(s))^*Qx^*, Qy^*]_H \, ds$$

Since by assumption $||S^{H}(\cdot)||_{\mathcal{L}(H)} \in L^{2}_{loc}[0,\infty)$, Hölder's inequality shows that Q_{t}^{H} extends to a bounded operator on H. The above identities show that this extension is self-adjoint.

By Lemma 1.10 we have $H_t \subset H$, which implies that for all $x^* \in E^*$ and $y^* \in E^*$,

$$[Q_t^H(Qx^*), Q_t^H(Qy^*)]_{\mathcal{H}_t} = [Q_t x^*, Q_t^H(Qy^*)]_{\mathcal{H}_t}$$

= $[Q_t x^*, Qy^*]_H$
= $\langle Q_t x^*, y^* \rangle$
= $[Q_t x^*, Q_t y^*]_{H_t}$
= $[Q_t^H(Qx^*), Q_t^H(Qy^*)]_{H_t}$

Hence the identity map restricted to $\operatorname{Im}(Q_t^H \circ Q) = \operatorname{Im} Q_t$ extends to an inner product preserving isomorphism of \mathcal{H}_t onto H_t .

The following examples illustrate the results of this section.

Example 1.12. Let $E = L^2[0, 1]$ and let w be the Wiener measure on E; thus, w is the centered Gaussian measure on E whose covariance operator Q is the integral operator on E defined by

$$(Qf)(s) = \int_0^1 (s \wedge \tau) f(\tau) \, d\tau.$$

The associated RKHS is the Hilbert space H of all absolutely continuous functions f on [0,1] for which f(0) = 0 and the a.e. derivative f' belongs to $L^2[0,1]$. The inner product of H is given by $[f,g]_H = [f',g']_E$.

Let \mathbf{S} be the nilpotent right shift semigroup on E, i.e.

$$S(t)f(s) = \begin{cases} f(s-t), & t \in [0,s];\\ 0, & \text{otherwise,} \end{cases} \quad s \in [0,1], \ t \ge 0.$$

We will show that $H_t = H_s$ for all t > 0 and s > 0. Since S(t) = 0 for $t \ge 1$ it is clear that $Q_t = Q_1$ and therefore $H_t = H_1$ for all $t \ge 1$. For this reason we will only consider $t \in (0, 1]$.

Denote by \mathbf{S}^{H} the restiction of \mathbf{S} to H and note that \mathbf{S}^{H} is a C_{0} -contraction semigroup on H. Therefore by Theorem 1.11, for all t > 0 we have $H_{t} = \mathcal{H}_{t}$ with identical norms.

We compute the space H_t explicitly. From

$$S^{H}(s)(S^{H}(s))^{*}h(\tau) = \chi_{[s,1]}(\tau)h(\tau), \qquad s \in [0,t], \ \tau \in [0,1], \ h \in H,$$

we have

$$Q_t^H h(\tau) = (t \wedge \tau) h(\tau), \qquad \tau \in [0, 1], \ h \in H.$$

Therefore,

$$H_t = \mathcal{H}_t = \operatorname{Im} \left((Q_t^H)^{1/2} \right)$$

= { $h \in H$: the function $\tau \mapsto (t \wedge \tau)^{-1/2} h(\tau)$ belongs to H }
= { $h \in H$: the function $\tau \mapsto \tau^{-1/2} h(\tau)$ belongs to H }.

Thus, H_t is independent of t, and its norm is given by

$$\|h\|_{H_t}^2 = \|h\|_{\mathcal{H}_t}^2 = [Q_t^H h, h]_H$$

= $[\chi_{[0,t]}h + (t \wedge \cdot)h', h']_E$
= $\int_0^t h(\tau)h'(\tau) d\tau + \int_0^1 (t \wedge \tau)(h'(\tau))^2 d\tau.$

From the representation of H_t it is clear that $S(s)H_t \subset H_t$ for all $s \ge 0$, so that a posteriori Theorem 1.9 applies. On the other hand, by control theoretic methods it is not difficult to show that the assumptions of Theorem 1.11 already imply the inclusion $S^H(t)H \subset \mathcal{H}_t$. Therefore by Lemma 1.10, $S(t)H_t \subset S(t)H \subset \mathcal{H}_t = H_t$ for all t > 0.

The next example shows that the inclusion $H_{t_0} \subset H_{t_1}$ may fail to be dense for certain $0 < t_0 < t_1$. The construction is based upon an example shown to the author by Szymon Peszat.

Example 1.13. Let $E = C_0[0, 1]$ be the Banach space of continuous real-valued functions f on [0, 1] with f(0) = 0. Let **S** be the nilpotent right shift semigroup on E. Fix $a \in (0, 1)$ arbitrary and let $Q \in \mathcal{L}(E^*, E)$ be the rank one operator defined by $Q\nu := \langle f_0, \nu \rangle f_0$, where $f_0 \in E$ is a function which is strictly positive on the interval (0, a) and vanishes on [a, 1]. Clearly Q is positive and symmetric. From

$$Q_t \nu = \int_0^t \langle f_0, S^*(s)\nu \rangle S(s) f_0 \, ds$$

it follows that for each t > 0 the RKHS H_t is contained in the closed linear span G_t of the set $\{S(s)f_0 : s \in [0, t]\}$.

Suppose $0 < t_0 < t_1 \leq 1$ are such that $t_1 - t_0 > a$. Then G_{t_0} , hence also H_{t_0} , is contained in the closed subspace E_{a+t_0} of E consisting of all functions vanishing on $[a + t_0, 1]$. On the other hand,

$$(Q_{t_1}\delta_{t_1})(t_1) = \left(\int_0^{t_1} S(s)Q\delta_{t_1-s} \, ds\right)(t_1)$$

= $\int_0^{t_1} f_0(t_1-s)(S(s)f_0)(t_1) \, ds$
= $\int_0^{t_1} (f_0(t_1-s))^2 \, ds > 0,$

where δ_{t_1} denotes the Dirac measure at t_1 . Since $t_1 > t_0 + a$, $Q_{t_1} \delta_{t_1} \notin E_{a+t_0}$. But if H_{t_0} were dense in H_{t_1} , we would have $Q_{t_1} \delta_{t_1} \in H_{t_1} = \overline{H_{t_0}}^{H_{t_1}} \subset \overline{H_{t_0}}^E \subset E_{a+t_0}$. Therefore the inclusion $H_{t_0} \subset H_{t_1}$ cannot be dense.

This example can be extended to show that Theorem 1.9 is the best possible:

Example 1.14. For each n let $E_n := C_0[0, 1]$, let \mathbf{S}_n be the nilpotent right shift semigroup on E_n , and let Q_n be as in Example 1.13 with $a_n := 1/n$. Let E be the c_0 -direct sum of the spaces E_n , and define \mathbf{S} and Q as direct sums of the \mathbf{S}_n and Q_n in the natural way. Then the inclusion $H_{t_0} \subset H_{t_1}$ fails to be dense for all $0 < t_0 < t_1 \leq 1$, this being the case in the kth summand whenever $t_1 - t_0 > 1/k$. On the other hand, the fact that S(t) = 0 for all $t \geq 1$ implies that $H_t = H_1$ for all $t \geq 1$.

Trivially, $S(1)H_{t_0} \subset H_{t_0}$ for all $t_0 > 0$. In particular this holds for any $t_0 \in (0, 1)$, although H_t is constant only after $t \ge 1$. This shows that Theorem 1.9 is the best possible in case $\max\{s, t_0\} = s$.

Similarly, for all s > 0 we have $S(s)H_1 \subset H_1$. In particular this holds for any $s \in (0, 1)$, although H_t is constant only after $t \ge 1$. This shows that Theorem 1.9 is also the best possible in case $\max\{s, t_0\} = t_0$.

2. The associated Gaussian measures μ_t

In this section and the next, E will always denote a separable real Banach space, and **S** is a fixed C_0 -semigroup on E.

It is not hard to see that for each positive symmetric $Q \in \mathcal{L}(E^*, E)$ there exists a unique finitely additive cylindrical measure μ , defined on the ring of all cylindrical sets of E, whose Fourier transform is given by

$$\hat{\mu}(x^*) = \exp\left(-\frac{1}{2}\langle Q \, x^*, x^*\rangle\right), \qquad x^* \in E^*.$$
(2.1)

In this section we fix a positive symmetric operator $Q \in \mathcal{L}(E^*, E)$ and make the following

Assumption 2.1. For each t > 0 the cylindrical measure μ_t associated with the positive symmetric operator $Q_t \in \mathcal{L}(E^*, E)$ is countably additive.

In other words, we assume that the operators Q_t are the covariances of centered Gaussian measures μ_t on the Borel σ -algebra of E.

Remark 2.2. We state two sufficient conditions for Assumption 2.1 to be satisfied:

- (i) E is a Hilbert space and Q is trace class (i.e. the cylindical measure associated with Q is countably additive);
- (ii) The cylindrical measure associated with Q is countably additive, $S(s)H \subset H$ for all $s \ge 0$, and

$$\int_0^t \|S(s)\|_{\mathcal{L}(H)}^2 \, ds < \infty$$

for all $t \ge 0$ [MS].

For the reader's convenience we reproduce the simple proofs; more information about (cylindrical) Gaussian measures can be found in the books [Ku], [VTC], and [DZ3].

(i): If (e_n) is an orthonormal basis in E, then by Fubini's theorem

$$\sum_{n=1}^{\infty} [Q_t e_n, e_n]_E = \int_0^t \sum_{n=1}^{\infty} [S(s)QS^*(s)e_n, e_n]_E \, ds$$
$$\leqslant \left(\sup_{0 \leqslant s \leqslant t} \|S(s)\|^2 \right) \cdot t \, \|Q\|_1 < \infty,$$

where $||Q||_1$ is the trace class norm of Q; we used the fact that for any bounded T, the operator TQT^* is trace class whenever Q is, with $||TQT^*||_1 \leq ||T|| ||Q||_1 ||T^*|| = ||T||^2 ||Q||_1$.

(ii): By Lemma 1.10, for each t > 0 there is a constant $K_t > 0$ such that

$$\langle Q_t x^*, x^* \rangle \leqslant K_t \langle Q x^*, x^* \rangle, \qquad \forall x^* \in E^*.$$

Countable additivity of μ_t then follows from [VTC, Corollary VI.3.4.2].

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On $B_b(E)$, the space of bounded, Borel measurable functions on E, the formula

$$(P(t)f)(x) := \int_E f(S(t)x + y) \, d\mu_t(y), \qquad x \in E,$$

defines a semigroup $\mathbf{P} = \{P(t)\}_{t \ge 0}$ of contractions. This semigroup will be referred to as the (non-symmetric) Ornstein-Uhlenbeck semigroup associated with \mathbf{S} and Q. In this section we state, without proof, a number of results about \mathbf{P} , the analogues of which are well-known in the Hilbert space setting. Their proofs carry over to the Banach space setting without difficulty and are therefore omitted.

Theorem 2.3. Let $x \in E$ and $t_0 > 0$ be fixed. The following assertions are equivalent:

- (i) $S(t_0)x \in H_{t_0};$
- (ii) For all $f \in B_b(E)$, the function $\varepsilon \mapsto P(t_0)f(\varepsilon x)$ is continuous at $\varepsilon = 0$;
- (iii) For all $f \in B_b(E)$ and $y \in E$, $P(t_0)f$ is smooth at y in the direction of x.

In this situation, the first directional derivative can be computed explicitly. For this purpose we introduce the following notation. If μ is a centered Gaussian measure on E, then $\phi^{\mu}: H \to L^2(E, \mu)$ denotes the isometric embedding uniquely defined by

$$\phi^{\mu}(Qx^*) := \langle x^*, \cdot \rangle,$$

where Q is the covariance operator of μ . For $h \in H$, the RHKS associated with Q, we will write ϕ_h^{μ} to denote the function $\phi^{\mu}(h) \in L^2(E,\mu)$. To see that this map is well-defined, recall that the support of μ is contained in the closure E_0 of H in E, whereas $Qx^* = Qy^*$ implies that $x^*|_{E_0} = x^*|_{E_0}$.

With this notation, the partial derivative $\partial P(t_0) f / \partial x$ is given by

$$\frac{\partial P(t_0)f}{\partial x}(y) = \int_E f(S(t_0)y + z)\phi_{S(t_0)x}^{\mu_{t_0}}(z) \, d\mu_{t_0}(z).$$

For the Wiener semigroup these results are due to Gross [Gr]; for E Hilbert they were extended to arbitrary semigroups **S** in [CG3].

The semigroup **P** is said to be strongly Feller at time $t_0 > 0$ if $P(t_0)f(\cdot)$ is a continuous function for all $f \in B_b(E)$. We refer to [DZ3] for more information in the Hilbert space setting.

Corollary 2.4. For $t_0 > 0$ fixed, the following conditions are equivalent:

- (i) **P** is strongly Feller at t_0 ;
- (ii) $S(t_0)E \subset H_{t_0}$.

Equivalence of the measures μ_t 3.

Two measures μ , ν are said to be *equivalent*, notation $\mu \sim \nu$, if they are absolutely continuous with respect to each other, i.e. if $\mu \ll \nu$ and $\nu \ll \mu$. We will study the question under what conditions we have equivalence $\mu_{t_0} \sim \mu_{t_1}$ for certain t_0 and t_1 . Our result is based on the following version of the Feldman-Hajek theorem, due to Tarieladze; see also the review paper [VT].

Theorem 3.1 [Ta]. Let μ , ν be two centered Gaussian measures on a Banach space E, and denote by $Q_{\mu}, Q_{\nu} \in \mathcal{L}(E^*, E)$, and H_{μ}, H_{ν} their covariance operators and RHKS's, respectively. Then $\mu \sim \nu$ if and only if the following two conditions are satisfied:

- (i) $H_{\mu} = H_{\nu};$
- (ii) $I j \circ V$ is Hilbert-Schmidt on H_{μ} , where $V : H_{\mu} \to H_{\nu}$ and $j : H_{\nu} \to H_{\mu}$ are defined by

$$VQ_{\mu}x^* := Q_{\nu}x^*, \qquad x^* \in E^*,$$

$$jh := h, \qquad h \in H_{\nu}.$$

Otherwise, $\mu \perp \nu$.

Throughout this section we consider a C_0 -semigroup **S** on E and a positive symmetric operator $Q \in \mathcal{L}(E^*, E)$ verifying Assumption 2.1.

Let $0 < t_0 < t_1 < \infty$. In terms of the operators $S_{t_1-t_0 \to t_1}(t_0) := S(t_0)|_{H_{t_1-t_0}} \in$ $\mathcal{L}(H_{t_1-t_0}, H_{t_1})$ (whose existence follows from Theorem 1.4) we can characterize equivalence of the measures μ_{t_0} and μ_{t_1} as follows:

Theorem 3.2. Let $0 < t_0 < t_1 < \infty$. Then $\mu_{t_0} \sim \mu_{t_1}$ if and only if the following two conditions are satisfied:

- (i) $||S_{t_1-t_0 \to t_1}(t_0)||_{\mathcal{L}(H_{t_1-t_0}, H_{t_1})} < 1;$ (ii) The operator $S_{t_1-t_0 \to t_1}(t_0)S^*_{t_1-t_0 \to t_1}(t_0)$ is Hilbert-Schmidt on H_{t_1} .

Proof: By Theorem 1.7, strict contractivity of $S_{t_1-t_0 \to t_1}(t_0)$ is equivalent to $H_{t_0} = H_{t_1}$. Next we consider the Hilbert-Schmidt condition. We have

$$Q_{t_1} - Q_{t_0} = S(t_0)Q_{t_1 - t_0}S^*(t_0) = S_{t_1 - t_0 \to t_1}(t_0)S^*_{t_1 - t_0 \to t_1}(t_0)Q_{t_1}.$$

Letting $j: H_{t_0} \to H_{t_1}$ be the identity map, it follows that $I - j \circ Q_{t_0} Q_{t_1}^{-1} : Q_{t_1} x^* \mapsto$ $Q_{t_1}x^* - Q_{t_0}x^*$ is Hilbert-Schmidt on H_{t_1} if and only if $S_{t_1-t_0 \to t_1}(t_0)S^*_{t_1-t_0 \to t_1}(t_0)$ is Hilbert-Schmidt on H_{t_1} .

Corollary 3.3. Suppose $0 < t_0 < t_1 < \infty$ are such that $\mu_{t_0} \sim \mu_{t_1}$.

- (i) For all $\delta \ge 0$ we have $\mu_{t_0+\delta} \sim \mu_{t_1+\delta}$;
- (ii) If $t_2 \in [t_1,\infty)$ is such that $\mu_t \sim \mu_{t_2}$ for all $t \in [t_2,\infty)$, then $\mu_t \sim \mu_{t_0}$ for all $t \in [t_0, \infty).$

Proof: (i): Fix $\delta \ge 0$ arbitrary. By Corollary 1.8, $H_{t_0+\delta} = H_{t_1+\delta}$. Hence from the identity

$$S_{(t_1+\delta)-(t_0+\delta)\to t_1+\delta}(t_0+\delta)S^*_{(t_1+\delta)-(t_0+\delta)\to t_1+\delta}(t_0)$$

= $S_{t_1\to t_1+\delta}(\delta) (S_{t_1-t_0\to t_1}(t_0)S^*_{t_1-t_0\to t_1}(t_0))S^*_{t_1\to t_1+\delta}(\delta)$

and Theorem 3.2 we conclude that $\mu_{t_0+\delta} \sim \mu_{t_1+\delta}$.

(ii): Pick $k \in \mathbb{N}$ such that $t_1 + k(t_1 - t_0) \ge t_2$. By (i) we have

$$\mu_{t_0} \sim \mu_{t_1} \sim \mu_{t_1 + (t_1 - t_0)} \sim \dots \sim \mu_{t_1 + k(t_1 - t_0)} \sim \mu_{t_2}$$

Hence, $\mu_t \sim \mu_{t_0}$ for all $t \in [t_2, \infty)$. But then, by another application of (i) we have, for all $t \in [t_0, t_2]$,

$$\mu_t = \mu_{t_0 + (t - t_0)} \sim \mu_{t_2 + (t - t_0)} \sim \mu_{t_0}.$$

It follows that $\mu_t \sim \mu_{t_0}$ for all $t \in [t_0, \infty)$.

If Q 'commutes' with S(s) for some s > 0, in the sense that $S(s)Q = QS^*(s)$, we can prove more. We start with a lemma (which was shown to the author by Ben de Pagter).

Lemma 3.4. Suppose μ and ν are centered Gaussian measures on E such that $H_{\mu} = H_{\nu}$. Let $V : H_{\mu} \to H_{\nu}$ and $j : H_{\nu} \to H_{\mu}$ be defined by

$$VQ_{\mu}x^* := Q_{\nu}x^*, \qquad x^* \in E^*,$$

$$jh := h, \qquad h \in H_{\nu}.$$

Then $V \circ j$ is positive and self-adjoint on H_{ν} , and $(V \circ j)^{\frac{1}{2}} \circ j^{-1}$ is an inner product preserving isomorphism of H_{μ} onto H_{ν} .

Proof: Clearly, the inner product $[\cdot, \cdot]_{H_{\mu}}$ defines a bounded symmetric bilinear form on H_{ν} . Consequently there exists a unique self-adjoint operator $V_1 \in \mathcal{L}(H_{\nu})$ such that

$$[jg, jh]_{H_{\mu}} = [V_1g, h]_{H_{\nu}}, \qquad \forall g, h \in H_{\nu}.$$

Moreover V_1 is positive and invertible. Let (e_n) be an orthonormal basis in H_{μ} . Then,

$$[V_1^{\frac{1}{2}}j^{-1}e_n, V_1^{\frac{1}{2}}j^{-1}e_m]_{H_{\nu}} = [V_1j^{-1}e_n, j^{-1}e_m]_{H_{\nu}} = [e_n, e_m]_{H_{\mu}} = \delta_{nm}.$$

Hence, $(V_1^{\frac{1}{2}}j^{-1}e_n)$ is an orthonormal basis for H_{ν} ; note that $V_1^{\frac{1}{2}} \in \mathcal{L}(H_{\nu})$ is surjective. Since $[V_1^{\frac{1}{2}}g, V_1^{\frac{1}{2}}h]_{H_{\nu}} = [jg, jh]_{H_{\mu}}$ for all $g, h \in H_{\nu}$, it follows that $V_1^{\frac{1}{2}} \circ j^{-1} : H_{\mu} \to H_{\nu}$ is an inner product preserving isometric isomorphism onto.

Returning to the map V, for all $x^*, y^* \in E^*$ we have

$$[Q_{\mu}x^*, Q_{\mu}y^*]_{H_{\mu}} = \langle Q_{\mu}y^*, x^* \rangle = [Q_{\nu}x^*, j^{-1}Q_{\mu}y^*]_{H_{\nu}} = [V(Q_{\mu}x^*), j^{-1}Q_{\mu}y^*]_{H_{\nu}}.$$

Hence via density, $[jg, jh]_{H_{\mu}} = [Vjg, h]_{H_{\nu}}$ for all $g, h \in H_{\nu}$. This shows that $V \circ j = V_1$.

Theorem 3.5. Suppose we have $S(s)Q = QS^*(s)$ for some s > 0. Let $t_0 > 0$ be fixed. If there exists $t_1 \in (t_0, \infty)$ such that $\mu_{t_1} \sim \mu_{t_0}$, then $\mu_t \sim \mu_{t_0}$ for all $t \in [t_0, \infty)$.

Proof: Step 1 - First assume $s = t_0$.

Fix $t_2 \ge 2t_0$ and assume for the moment that $\mu_{t_2} \sim \mu_{t_0}$. We will prove that $\mu_t \sim \mu_{t_2}$ for all $t \in [t_2, \infty)$.

Clearly, $H_{t_2} = H_{t_0}$. Fix $t \ge t_2$. In view of $t_2 - t_0 \ge t_0$, by Corollary 1.8 we also have $H_{t-t_0} = H_{t_2-t_0} = H_{t_0}$ and $H_t = H_{t_0}$. Define $V : H_{t_2-t_0} \to H_{t-t_0}$ by

$$V: Q_{t_2-t_0}x^* \mapsto Q_{t-t_0}x^*, \qquad x^* \in E^*.$$

From $S(t_0)Q = QS^*(t_0)$ we have $S(t_0)Q_\tau = Q_\tau S^*(t_0)$ for all $\tau > 0$ and hence

$$S(t_0)VQ_{t_2-t_0}x^* = S(t_0)Q_{t-t_0}x^* = Q_{t-t_0}S^*(t_0)x^*$$
$$= VQ_{t_2-t_0}S^*(t_0)x^* = VS(t_0)Q_{t_2-t_0}x^*, \qquad x^* \in E^*.$$

Letting $j: H_{t-t_0} \to H_{t_2-t_0}$ be the identity operator, it follows that on H_{t-t_0} we have

$$S_{t-t_0 \to t-t_0}(t_0) \circ (V \circ j) = (V \circ j) \circ S_{t-t_0 \to t-t_0}(t_0).$$

By Lemma 3.4, $V \circ j$ is positive on H_{t-t_0} , and $U := (V \circ j)^{\frac{1}{2}} \circ j^{-1}$ is an inner product preserving isomorphism of $H_{t_2-t_0}$ onto H_{t-t_0} . Moreover, by functional calculus we have

$$S_{t-t_0 \to t-t_0}(t_0) \circ (V \circ j)^{\frac{1}{2}} = (V \circ j)^{\frac{1}{2}} \circ S_{t-t_0 \to t-t_0}(t_0).$$

Multiplying on the right with j^{-1} gives

$$S_{t-t_0 \to t-t_0}(t_0) = (V \circ j)^{\frac{1}{2}} \circ S_{t-t_0 \to t-t_0}(t_0) \circ j^{-1} \circ U^* = U \circ S_{t_2-t_0 \to t_2-t_0}(t_0) \circ U^*.$$

Therefore,

$$S_{t-t_0 \to t}(t_0) = j_{t-t_0 \to t} \circ U \circ j_{t_2 \to t_2 - t_0} \circ S_{t_2 - t_0 \to t_2}(t_0) \circ U^*,$$

where $j_{t-t_0 \to t}$ and $j_{t_2 \to t_2 - t_0}$ are the identity maps from H_{t-t_0} to H_t and from H_{t_2} to $H_{t_2-t_0}$, respectively. Using the equivalence $\mu_{t_2} \sim \mu_{t_0}$ and Theorem 3.2, we conclude that

$$\begin{split} S_{t-t_0 \to t}(t_0) S_{t-t_0 \to t}^*(t_0) \\ &= (j_{t-t_0 \to t} \circ U \circ j_{t_2 \to t_2 - t_0} \circ S_{t_2 - t_0 \to t_2}(t_0) \circ U^*) \\ &\quad \circ \left(U \circ S_{t_2 - t_0 \to t_2}^*(t_0) \circ j_{t_2 \to t_2 - t_0}^* \circ U^* \circ j_{t-t_0 \to t}^* \right) \\ &= j_{t-t_0 \to t} \circ U \circ j_{t_2 \to t_2 - t_0} \circ \left(S_{t_2 - t_0 \to t_2}(t_0) S_{t_2 - t_0 \to t_2}^*(t_0) \right) \circ j_{t_2 \to t_2 - t_0}^* \circ U^* \circ j_{t-t_0 \to t}^* \end{split}$$

is Hilbert-Schmidt on H_t . By another application of Theorem 3.2 it follows that $\mu_t \sim \mu_{t_0}$.

Step 2 - Using Step 1, we now prove the theorem for the case $s = t_0$.

Let $k \in \mathbb{N}$ be any integer such that $t_1 + k(t_1 - t_0) \ge 2t_0$. By Corollary 3.3 (i) we have

$$\mu_{t_0} \sim \mu_{t_1} \sim \mu_{t_1+(t_1-t_0)} \sim \dots \sim \mu_{t_1+k(t_1-t_0)}$$

We may apply Step 1 to $t_2 := t_1 + k(t_1 - t_0)$. It follows that $\mu_t \sim \mu_{t_2}$ for all $t \ge t_2$. But then Corollary 3.3 (ii) shows that $\mu_t \sim \mu_{t_0}$ for all $t \in [t_0, \infty)$.

Step 3 - We now prove the general case. Choose an integer m such that $ms \ge t_1$. Then $S(ms)Q = QS^*(ms)$; further $\mu_{t_0} \sim \mu_{t_1}$ implies $\mu_{ms} \sim \mu_{t_1+ms-t_0}$ by Corollary 3.3 (i). Hence we may apply Step 2 to $\tau_0 := ms$ and $\tau_1 := t_1 + ms - t_0$. It follows that $\mu_t \sim \mu_{ms}$ for all $t \in [ms, \infty)$. But then we apply Corrollary 3.3 (ii) to t_0, t_1 , and $t_2 := ms$ to see that $\mu_t \sim \mu_{t_0}$ for all $t \in [t_0, \infty)$.

In view of this result and by the analogy to Corollary 1.8 we conjecture that $\mu_{t_1} \sim \mu_{t_0}$ always implies $\mu_t \sim \mu_{t_0}$ for all $t \in [t_0, \infty)$.

The following corollary gives necessary and sufficient conditions for the situation described by Theorem 3.5:

Corollary 3.6. Suppose we have $S(s)Q = QS^*(s)$ for some s > 0. Then the following assertions are equivalent:

(i) $\mu_t \sim \mu_{t_0}$ for all $t \in [t_0, \infty)$;

(ii) $S(t_0)H_{t_0} \subset H_{t_0}$ and $S(t_0)|_{H_{t_0}}(S(t_0)|_{H_{t_0}})^*$ is Hilbert-Schmidt on H_{t_0} .

Proof: Assume (i). Then in particular $\mu_{2t_0} \sim \mu_{t_0}$, and Theorem 3.2 shows that $S_{t_0 \to 2t_0}(t_0)(S_{t_0 \to 2t_0}(t_0))^*$ is Hilbert-Schmidt on H_{2t_0} . But then

$$S_{t_0 \to t_0}(t_0)(S_{t_0 \to t_0}(t_0))^* = (j_{2t_0 \to t_0} \circ S_{t_0 \to 2t_0}(t_0)) \circ (S_{t_0 \to t_0}^*(t_0) \circ j_{2t_0 \to t_0}^*)$$

is Hilbert-Schmidt on H_{t_0} . This gives (ii).

Conversely, if (ii) holds, then

$$S_{t_0 \to 2t_0}(t_0)(S_{t_0 \to 2t_0}(t_0))^* = (j_{t_0 \to 2t_0} \circ S_{t_0 \to t_0}(t_0)) \circ (S_{t_0 \to t_0}^*(t_0) \circ j_{t_0 \to 2t_0}^*)$$
$$= j_{t_0 \to 2t_0} \circ S_{t_0 \to t_0}(2t_0) \circ j_{t_0 \to 2t_0}^*$$

Hilbert-Schmidt on H_{2t_0} . Theorem 3.2 shows that $\mu_{2t_0} \sim \mu_{t_0}$, and therefore (i) holds by Theorem 3.5.

Under the assumptions that E is Hilbert and that for all s > 0 the operator S(s) is self-adjoint on E and commutes with Q, this result is essentially equivalent to [NZ, Theorem 3.1].

Remark 3.7. If $S(t_0)Q = QS^*(t_0)$, then (ii) may be replaced by (ii)' $S(t_0)H_{t_0} \subset H_{t_0}$ and the restriction $S(2t_0)|_{H_{t_0}}$ is Hilbert-Schmidt on H_{t_0} . This follows from (ii) once we show that the restriction $S(t_0)|_{H_{t_0}}$ is self-adjoint on H_{t_0} :

$$\begin{split} [(S(t_0)|_{H_{t_0}})^* Q_{t_0} x^*, Q_{t_0} y^*]_{H_{t_0}} &= [Q_{t_0} x^*, S(t_0)|_{H_{t_0}} Q_{t_0} y^*]_{H_{t_0}} \\ &= [Q_{t_0} x^*, Q_{t_0} S^*(t_0) y^*]_{H_{t_0}} \\ &= [Q_{t_0} x^*, (S(t_0)|_{H_{t_0}})^* Q_{t_0} y^*]_{H_{t_0}} \\ &= [S(t_0)|_{H_{t_0}} Q_{t_0} x^*, Q_{t_0} y^*]_{H_{t_0}} \end{split}$$

for all $x^*, y^* \in E^*$.

As the following two examples show, it may happen that $H_t = H_{t_0}$ for all $t \in [t_0, \infty)$ although $\mu_t \perp \mu_s$ for all $t \neq s \in [t_0, \infty)$. A third example is given in Section 6 below.

Example 3.8. Let E be an infinite-dimensional Hilbert space and let μ be a non-degenerate centered Gaussian measure on E with covariance operator Q. Let **S** be a periodic C_0 -semigroup on E, with period 1. Then Assumption 2.1 holds, we have $Q_k = kQ_1$ for all k = 1, 2, ..., and consequently $H_k = H_1$ for all such k. Hence, $H_t = H_1$ for all $t \in [1, \infty)$. On the other hand, let us suppose that $\mu_t \sim \mu_s$ for certain $t, s \in [1, \infty)$ with t < s, then for any integer $k \ge t$ we also have $\mu_k \sim \mu_{s+k-t}$ by Corollary 3.3. But S(k) = I commutes with Q, and therefore Theorem 3.5 implies $\mu_{\tau} \sim \mu_k$ for all $\tau \in [k, \infty)$; in particular, $\mu_k \sim \mu_l$ for all integers l > k. But these measures have covariances kQ and lQ, respectively, and therefore they are singular by the Feldman-Hajek theorem; a contradiction.

Example 3.9. We continue Example 1.12. By Remark 2.2 (i), each of the operators Q_t is the covariance of a centered Gaussian measure μ_t on $E = L^2[0, 1]$. We will show that $\mu_t \perp \mu_s$ if $t \in (0, 1)$ and $s \neq t$, whereas it is trivial that $\mu_t = \mu_s$ whenever $t \ge 1$ and $s \ge 1$.

Fix $t \in (0,1)$ and s > 0, $s \neq t$. Since $\mu_s = \mu_1$ if $s \ge 1$ we may assume that $s \in (0,1]$. By interchanging the roles of t and s if necessary, we may also assume that t < s, say s = t + h for some $h \in (0, 1 - t]$. Let F denote the closed subspace of H_{t+h} consisting of all functions with support in [t,1]. For all $f \in F$, $S_{h\to t+h}(t)S_{h\to t+h}^*(t)f = f$, so $S^H(t)(S^H(t))^*|_F = I_F$, the identity operator on F. Since dim $F = \infty$ it follows that $S_{h\to t+h}(t)S_{h\to t+h}^*(t)$ is not compact on H_{t+h} and therefore not Hilbert-Schmidt. This shows that $\mu_t \perp \mu_{t+h} = \mu_s$.

4. Computation of the Radon-Nikodym derivative

It is possible to give an explicit expression for the Radon-Nikodym density $d\mu_{t_1}/d\mu_{t_0}$ whenever we have $\mu_{t_0} \sim \mu_{t_1}$. This will occupy us in the present section.

We start by recalling some notation and results concerning second quantization. For more details we refer to [Ne] and the book [Si]. Fix a centered Gaussian measure μ on E with covariance operator $Q \in \mathcal{L}(E^*, E)$, and let H denote the associated RKHS. Let $\phi^{\mu} : H \mapsto L^2(E, \mu)$ be the isometric embedding from H into $L^2(E, \mu)$ defined by $\phi^{\mu}(Qx^*) = \langle x^*, \cdot \rangle$ as in Section 2. Whenever the measure μ is understood, we omit it from the notation and write ϕ_h to denote the function $\phi(h) = \phi^{\mu}(h) \in L^2(E, \mu)$.

Let $(H_n)_{n \in \mathbb{N}}$ be the sequence of Hermite polynomials and denote by \mathcal{H}_n the closure in $L^2(E,\mu)$ of the linear span of the set $\{H_n(\phi_h): \|h\|_H = 1\}$. Note that \mathcal{H}_0 is the one-dimensional subspace spanned by the constant one function and that \mathcal{H}_1 is precisely the image of H under the isometry ϕ . One has the orthogonal Wiener-Itô decomposition,

$$L^2(E,\mu) = \bigoplus_{n \in \mathbb{N}} \mathcal{H}_n.$$

The orthogonal projection onto \mathcal{H}_n will be denoted by I_n .

For all $h \in H$, the functions

$$K_h(x) := \exp\left(\phi_h(x) - \frac{1}{2} \|h\|_H^2\right),$$

belong to $L^2(E,\mu)$, their linear span is dense in $L^2(E,\mu)$, and we have the identity

$$K_h = \sum_{n=0}^{\infty} \frac{1}{n!} I_n(\phi_h^n), \qquad h \in H$$

Now assume we have two pairs (E_0, μ_0) and (E_1, μ_1) , and let $T \in \mathcal{L}(H_0, H_1)$ be a contraction. The second quantization of T is the contraction $\Gamma(T) \in \mathcal{L}(L^2(E_0, \mu_0), L^2(E_1, \mu_1))$ defined by

$$\Gamma(T)\big(I_n(\phi_{h_1}^{k_1}\cdot\ldots\cdot\phi_{h_j}^{k_j})\big):=I_n(\phi_{Th_1}^{k_1}\cdot\ldots\cdot\phi_{Th_j}^{k_j}),$$

where it is assumed that $k_1 + \ldots + k_j = n$.

Now let a positive symmetric operator $Q \in \mathcal{L}(E^*, E)$ and a C_0 -semigroup **S** on E be given such that Assumption 2.1 holds. Let **P** be the Ornstein-Uhlenbeck semigroup on $B_b(E)$ associated with **S** and Q. We are going to apply second quantization to $E_0 = E_1 = E, \ \mu_0 := \mu_{t_0+h}, \ \mu_1 := \mu_h$, and the adjoint $S^*_{h \to t_0+h}(t_0) \in \mathcal{L}(H_{t_0+h}, H_h)$ of the Hilbert space contraction $S_{h \to t_0+h}(t_0) \in \mathcal{L}(H_h, H_{t_0+h})$.

Theorem 4.1. For all $t_0 > 0$ and h > 0, the operator $P(t_0)$ extends to a contraction from $L^2(E, \mu_{t_0+h})$ into $L^2(E, \mu_h)$. This extension is realized as the second quantization of $S^*_{h \to t_0+h}(t_0)$:

$$P(t_0) = \Gamma(S^*_{h \to t_0 + h}(t_0)).$$

Proof: We denote the image measure of μ_t with respect to an element $x^* \in E^*$ by $\langle x^*, \mu_t \rangle$. For all $x^* \in E^*$ we then have

$$\begin{split} P(t_0)K_{Q_{t_0+h}x^*}(x) &= \int_E \exp\left(\langle x^*, S(t_0)x + y \rangle - \frac{1}{2} \|Q_{t_0+h}x^*\|_{H_{t_0+h}}^2\right) d\mu_{t_0}(y) \\ &= K_{Q_{t_0+h}x^*}(S(t_0)x) \int_E \exp\left(\langle x^*, y \rangle\right) d\mu_{t_0}(y) \\ &= K_{Q_{t_0+h}x^*}(S(t_0)x) \int_{\mathbb{R}} \exp(s) d\langle x^*, \mu_{t_0} \rangle(s) \\ &= K_{Q_{t_0+h}x^*}(S(t_0)x) \exp\left(\frac{1}{2} \|Q_{t_0}x^*\|_{H_{t_0}}^2\right) \\ &= \exp\left(\langle x^*, S(t_0)x \rangle - \frac{1}{2} \left(\|Q_{t_0+h}x^*\|_{H_{t_0+h}}^2 - \|Q_{t_0}x^*\|_{H_{t_0}}^2\right)\right) \\ &= \exp\left(\langle x^*, S(t_0)x \rangle - \frac{1}{2} \|Q_hS^*(t_0)x^*\|_{H_h}^2\right) \\ &= K_{Q_hS^*(t_0)x^*}(x) \\ &= K_{S_{h\to t_0+h}^*(t_0)Q_{t_0+h}x^*}(x). \end{split}$$

Hence the identity

$$K_g = \sum_{n=0}^{\infty} \frac{1}{n!} I_n(\phi_g^n), \qquad g \in H_{t_0+h},$$

implies

$$P(t_0)K_{Q_{t_0+h}x^*} = K_{S_{h\to t_0+h}^*(h)Q_{t_0+h}x^*}$$
$$= \sum_{n=0}^{\infty} \frac{1}{n!} I_n(\phi_{S_{h\to t_0+h}^*(t_0)Q_{t_0+h}x^*})$$
$$= \Gamma(S_{h\to t_0+h}^*(t_0))K_{Q_{t_0+h}x^*}.$$

By a density argument, it follows that

$$P(t_0)K_g = \Gamma(S^*_{h \to t_0 + h}(t_0))K_g, \qquad \forall g \in H_{t_0 + h}.$$

Since the linear span of the functions K_g , $g \in H_{t_0+h}$, is dense in $L^2(E, \mu_{t_0+h})$, this proves the theorem.

The next aim is to apply the so-called *Mehler formula* for second quantized operators to the above situation.

To this end, we consider the situation of two pairs (E_0, μ_0) and (E_1, μ_1) , with μ_k a centered Gaussian measure on E_k with RHKS H_k ; k = 0, 1. The following result, due to Feyel and La Pradelle, shows that every bounded operator in $\mathcal{L}(H_0, H_1)$ has an extension to a linear μ_0 -measurable extension from E_0 into E_1 . Recall that a mapping $f : E_0 \to E_1$ is μ_0 -measurable if $f^{-1}(B)$ belongs to the μ_0 -completion of the Borel σ -algebra of E_0 , for all Borel sets $B \subset E_1$.

Proposition 4.2 [F-LP, Théorème 5]. Let $T \in \mathcal{L}(H_0, H_1)$. Then there exists a μ_0 -measurable linear operator \overline{T} from E_0 into E_1 which extends T. This extension is μ_0 -essentially unique in the sense that any two such μ_0 -measurable linear extensions agree μ_0 -a.e. Moreover, for all $h \in H_1$ we have $\phi_h^{\mu_1}(\overline{T}(x)) = \phi_{T^*h}^{\mu_0}(x)$ for μ_0 -almost all $x \in E_0$.

The uniqueness part implies that for a bounded operator $T \in \mathcal{L}(E_0, E_1)$ which maps H_0 into H_1 we have $T = \overline{T|_{H_0}} \mu_0$ -a.e.

In terms of these extensions, one has the following *Mehler formula* for the second quantization of a Hilbert space contraction:

Proposition 4.3 [F-LP, Théorème 10]. Let T be a contraction in $\mathcal{L}(H_0, H_1)$. Then for all $f \in L^2(E_0, \mu_0)$ and μ_1 -almost all $x \in E_1$ we have

$$\Gamma(T)f(x) = \int_{E_0} f\left(\overline{T^*}(x) + \overline{\sqrt{I - T^*T}}(y)\right) d\mu_0(y).$$

This result motivates the consideration of the image measure of μ_0 under the μ_0 -measurable transformation $\sqrt{I - T^*T}$. Let us denote this measure by μ_0^T . One has the following extension of a result of Shale [Sh]:

Proposition 4.4 [F-LP, Proposition 12]. Let $T \in \mathcal{L}(H_0, H_1)$ be a strict contraction such that T^*T is Hilbert-Schmidt on H_0 . Then $\mu_0^T \ll \mu_0$, with Radon-Nikodym derivative μ_0 -a.e. given by

$$\frac{d\mu_0^T}{d\mu_0}(x) = \frac{1}{\sqrt{\det(I - T^*T)}} \exp\left(-\frac{1}{2} \left\| \overline{\left(\sqrt{I - T^*T}\right)^{-1} (T^*T)^{\frac{1}{2}}}(x) \right\|^2\right).$$

If $\mu_{t_0} \sim \mu_{t_0+h}$ for some $t_0 > 0$ and h > 0, then $S_{h \to t_0+h}(t_0)S^*_{h \to t_0+h}(t_0)$ is Hilbert-Schmidt on H_{t_0+h} . If we assume that $S_{h \to t_0+h}(t_0)$ itself is Hilbert-Schmidt as an operator from H_h into H_{t_0+h} we can prove more:

Theorem 4.5. Suppose we have $\mu_{t_0} \sim \mu_{t_0+h}$ for some $t_0 > 0$ and h > 0. If $S_{h\to t_0+h}(t_0)$ is Hilbert-Schmidt from H_h into H_{t_0+h} , then the Radon-Nikodym derivative $d\mu_{t_0}/d\mu_{t_0+h}(x)$ is μ_{t_0+h} -a.e. given by

$$\frac{1}{\sqrt{\det(I-T)}} \exp\left(-\frac{1}{2} \left\|\overline{\left(\sqrt{I-T}\right)^{-1} T^{\frac{1}{2}}}(x)\right\|^{2}\right),$$

where $T := S_{h \to t_0 + h}(t_0) S^*_{h \to t_0 + h}(t_0).$

Proof: We note that for all $0 \leq f \in B_b(E)$ we have

$$P(t_0)f(x) = \int_E f(S(t_0)x + y) \, d\mu_{t_0}(y) = \int_E f(S(t_0)x + y) \frac{d\mu_{t_0}}{d\mu_{t_0+h}}(y) \, d\mu_{t_0+h}(y).$$

Combining this with Theorem 4.1 and Proposition 4.3, we see that

$$P(t_0)f(x) = \Gamma(S_{h \to t_0+h}^*(t_0))f(x)$$

= $\int_E f\left(\overline{S_{h \to t_0+h}(t_0)}(x) + \overline{\sqrt{I-T}}(y)\right) d\mu_{t_0+h}(y)$
= $\int_E f\left(S(t_0)x + \overline{\sqrt{I-T}}(y)\right) d\mu_{t_0+h}(y)$

By Proposition 4.4, the image measure of μ_{t_0+h} under the μ_{t_0+h} -measurable transformation $\sqrt{I-T}$ is absolutely continuous with respect to μ_{t_0+h} , with Radon-Nikodym derivative given, for μ_{t_0+h} -a.a. $y \in E$, by

$$\frac{1}{\sqrt{\det(I-T)}} \exp\left(-\frac{1}{2} \left\|\overline{\left(\sqrt{I-T}\right)^{-1} T^{\frac{1}{2}}}(y)\right\|^{2}\right).$$

Hence,

$$P(t_0)f(x) = \frac{1}{\sqrt{\det(I-T)}} \int_E f(S(t_0)x + y) \exp\left(-\frac{1}{2} \left\| \overline{\left(\sqrt{I-T}\right)^{-1} T^{\frac{1}{2}}}(y) \right\|^2 \right) d\mu_{t_0+h}(y),$$

and the desired result follows by comparing the two identities for $P(t_0)f(x)$.

5. Continuous dependence of the Radon-Nikodym derivative

Thoughout this section E is a separable real Banach space, $Q \in \mathcal{L}(E^*, E)$ is positive and symmetric, and **S** is a C_0 -semigroup on E such that Assumption 2.1 is verified. We will show that for t_1 fixed the Radon-Nikodym derivative $d\mu_{t_0}/d\mu_{t_1}$ depends continuously upon t_0 .

Our first aim is to establish a result concerning continuity of determinants.

Lemma 5.1. Fix $\tau > 0$. For all $g \in H_{\tau}$ we have

$$\lim_{h \downarrow 0} \|S_{\tau-h \to \tau}(h)S^*_{\tau-h \to \tau}(h)g - g\|_{H_{\tau}} = 0.$$

Proof: Fix $x^* \in E^*$ and $h \in [0, \tau)$. Writing $T_h := S_{\tau - h \to \tau}(h)$, for all $y^* \in E^*$ we have

$$[T_h^* Q_\tau x^*, T_h^* Q_\tau y^*]_{H_\tau} = [Q_{\tau-h} S^*(h) x^*, Q_{\tau-h} S^*(h) y^*]_{H_{\tau-h}},$$

= $\langle S(h) Q_{\tau-h} S^*(h) x^*, y^* \rangle$
= $\langle Q_\tau x^* - Q_h x^*, y^* \rangle$
= $[Q_\tau x^* - Q_h x^*, Q_\tau y^*]_{H_\tau}.$

Hence,

$$[T_h T_h^* Q_\tau x^* - Q_\tau x^*, Q_\tau y^*]_{H_\tau} = -[Q_h x^*, Q_\tau y^*]_{H_\tau}$$

Taking the supremum with respect to all $Q_{\tau}y^*$ of norm ≤ 1 , it follows that

$$||T_h T_h^* Q_\tau x^* - Q_\tau x^*||_{H_\tau} = ||Q_h x^*||_{H_\tau} \leq ||Q_h x^*||_{H_h},$$

the inequality being a consequence of Corollary 1.5. As $h \downarrow 0$ the right hand side tends to 0. Since $||T_h|| \leq 1$ for all h by Theorem 1.4, the lemma now follows by a density argument.

Lemma 5.2. Let H_0 and H_1 be separable Hilbert spaces, let $S \in \mathcal{L}(H_0, H_1)$ be Hilbert-Schmidt and let $(T_n) \subset \mathcal{L}(H_0)$ be a sequence of operators converging to I strongly. Then

$$\lim_{n \to \infty} ST_n S^* = SS^*$$

in the space $\mathcal{L}_1(H_1)$ of trace class operators on H_1 .

Proof: The lemma is obvious if S is a rank one operator. By taking linear combinations, it also holds for finite rank operators S. The general case then follows from a 3ε -argument, approximating S in the Hilbert-Schmidt norm by finite rank operators.

The preceding two lemmas combined with the fact [GGK, p. 119] that the mapping $T \mapsto \det(I-T)$ is continuous with respect to the trace class norm lead to the following result:

Lemma 5.3. Let $0 < t_0 < t_1$ be fixed and assume that the operator $S_{t_1-t_0 \to t_1}(t_0) \in$ $\mathcal{L}(H_{t_1-t_0}, H_{t_1})$ is Hilbert-Schmidt. Then the function

$$h \mapsto \det(I - S_{t_1 - t_0 - h \to t_1}(t_0 + h)S^*_{t_1 - t_0 - h \to t_1}(t_0 + h)), \qquad h \in [0, t_1 - t_0).$$

is continuous.

In the following lemma, $C_b(\Omega)$ denotes the space of bounded real-valued continuous functions on a topological space Ω .

Suppose $\tilde{f} \in C_b(\mathbb{R}^n)$ and $x_1^*, ..., x_n^* \in E^*$ are given, and define $f \in$ Lemma 5.4. $C_b(E)$ by f

$$f(x) := f(\langle x_1^*, x \rangle, ..., \langle x_n^*, x \rangle), \qquad x \in E.$$

Then for all $t_0 \ge 0$ and $x \in E$ we have we have

$$\lim_{h \downarrow 0} P(t_0 + h)f(x) - P(t_0)f(x) = 0.$$

Proof: We have

$$\begin{split} P(t_0 + h)f(x) &= \int_E \tilde{f}(\langle x_1^*, S(t_0 + h)x + y \rangle, ..., \langle x_n^*, S(t_0 + h)x + y \rangle) \, d\mu_{t_0 + h}(y) \\ &= \int_E \tilde{f}(\langle x_1^*, z \rangle, ..., \langle x_n^*, z \rangle) \, d\mu_{t_0 + h}^{(S(t_0 + h)x)}(z) \\ &= \int_{\mathbb{R}^n} \tilde{f}(\tau_1, ..., \tau_n) \, d\nu_{t_0 + h}^{(S(t_0 + h)x)}(\tau), \end{split}$$

where $\mu_{t_0+h}^{(S(t_0+h)x)}$ is the translation of μ_{t_0+h} along $S(t_0+h)x$, and $\nu_{t_0+h}^{(S(t_0+h)x)}$ is the image measure on \mathbb{R}^n of $\mu_{t_0+h}^{(S(t_0+h)x)}$ under the map $T: E \to \mathbb{R}^n$ given by Tz := $(\langle x_1^*, z \rangle, ..., \langle x_n^*, z \rangle)$. Thus, the Gaussian measure $\nu_{t_0+h}^{(S(t_0+h)x)}$ has mean $(\langle x_1^*, S(t_0+h)x \rangle, ..., \langle x_n^*, S(t_0+h)x \rangle)$ and covariance $TQ_{t_0+h}T^*$. By Lévy's theorem,

$$\lim_{h \downarrow 0} \nu_{t_0+h}^{(S(t_0+h)x)} = \nu_{t_0}^{(S(t_0)x)} \quad \text{weakly}$$

But then

$$\lim_{h \downarrow 0} P(t_0 + h) f(x) = \lim_{h \downarrow 0} \int_{\mathbb{R}^n} \tilde{f}(\tau_1, ..., \tau_n) \, d\nu_{t_0 + h}^{(S(t_0 + h)x)}(\tau)$$
$$= \int_{\mathbb{R}^n} \tilde{f}(\tau_1, ..., \tau_n) \, d\nu_{t_0}^{(S(t_0)x)}(\tau)$$
$$= P(t_0) f(x).$$

It is well-known that the space of all cylindrical $C_b(E)$ -functions as considered in Lemma 5.4 are dense in $L^2(E,\mu)$, for any Gaussian measure μ defined on the Borel σ -algebra of E. This will be used in the following theorem, which is the main result of this section.

Theorem 5.5. Assume that $\mu_t \sim \mu_{t_0}$ for all $t \in [t_0, \infty)$ and that for all h > 0the operator $S_{h \to t_0+h}(t_0)$ is Hilbert-Schmidt from H_h to H_{t_0+h} . Fix $t_1 > t_0$, and for $\tau \in [t_0, t_1]$ let $g_\tau := d\mu_\tau/d\mu_{t_1}$ denote the Radon-Nikodym derivative. Then

$$\lim_{h \downarrow 0} \|g_{t_0+h} - g_{t_0}\|_{L^2(E,\mu_{t_1})} = 0.$$

Proof: The proof is divided into two steps.

Step 1 - We first prove that

$$\lim_{h\downarrow 0} \|g_{t_0+h}\|_{L^2(E,\mu_{t_1})} = \|g_{t_0}\|_{L^2(E,\mu_{t_1})}$$

For $\tau \in [t_0, t_1]$, we define $T_{\tau} \in \mathcal{L}(H_{t_1})$ by

$$T_{\tau} := S_{t_1 - \tau \to t_1}(\tau) S^*_{t_1 - \tau \to t_1}(\tau).$$

Then,

$$\begin{split} \|g_{t_0+h}\|_{L^2(E,\mu_{t_1})}^2 &= \\ &= \frac{1}{\det(I - T_{t_0+h})} \int_E \exp\left(-\left\|\overline{\left(\sqrt{I - T_{t_0+h}}\right)^{-1} \sqrt{T_{t_0+h}}(x)}\right\|^2\right) d\mu_{t_1}(x) \\ &\leqslant \frac{1}{\det(I - T_{t_0+h})} \int_E \exp\left(-\frac{1}{2} \left\|\overline{\left(\sqrt{I - T_{t_0+h}}\right)^{-1} \sqrt{T_{t_0+h}}(x)}\right\|^2\right) d\mu_{t_1}(x) \\ &= \frac{1}{\sqrt{\det(I - T_{t_0+h})}} \|g_{t_0+h}\|_{L^1(E,\mu_{t_1})} \\ &= \frac{1}{\sqrt{\det(I - T_{t_0+h})}}. \end{split}$$

Therefore Step 1 is a consequence of Lemma 5.3.

Step 2 - The cylindrical functions as described in Lemma 5.4 are dense in $L^2(E, \mu_{t_1})$, and for each such f we have

$$\lim_{h \downarrow 0} \int_{E} f(x)(g_{t_0+h}(x) - g_{t_0}(x)) \, d\mu_{t_1}(x) = \lim_{h \downarrow 0} P(t_0+h)f(0) - P(t_0)f(0) = 0$$

Since by Step 1 the norms $||g_{t_0+h}||_{L^2(E,\mu_{t_1})}$ remain bounded as $h \downarrow 0$, it follows that $\lim_{h\downarrow 0} g_{t_0+h} = g_{t_0}$ weakly in $L^2(E,\mu_{t_1})$. Together with Step 1 this implies that $\lim_{h\downarrow 0} g_{t_0+h} = g_{t_0}$ strongly in $L^2(E,\mu_{t_1})$.

We will apply this result to show that under certain conditions the Ornstein-Uhlenbeck semigroup **P** associated with **S** and *Q* is pointwise continuous for $t \ge t_0$, uniformly on bounded sets in *E*, in the space BUC(E) of bounded real-valued uniformly continuous functions on *E*. Before doing so we make the following simple observation.

Proposition 5.6. Let $f \in C_b(E)$ and $t_0 > 0$ be fixed. If $S(t_0)$ is compact on E and

$$\lim_{h \downarrow 0} \left(\sup_{x \in K} |P(h)f(x) - f(x)| \right) = 0$$
(5.1)

for all compact sets $K \subset E$, then for all bounded sets $B \subset E$ we have

$$\lim_{h \downarrow 0} \left(\sup_{x \in B} \left| P(t_0 + h) f(x) - P(t_0) f(x) \right| \right) = 0$$

Proof: Given $\varepsilon > 0$ and a bounded set $B \subset E$, let $K_0 := \overline{S(t_0)B}$ and let $K_1 \subset E$ be a compact set such that $\mu_{t_0}(K_1) > 1 - \varepsilon$. Writing $g_h := P(h)f - f$, we have $\lim_{h \downarrow 0} g_h = 0$ uniformly on the compact set $\{y_0 + y_1 : y_0 \in K_0, y_1 \in K_1\}$, and hence

$$\begin{split} \lim_{h \downarrow 0} \left(\sup_{x \in B} |P(t_0 + h)f(x) - P(t_0)f(x)| \right) \\ &= \lim_{h \downarrow 0} \left(\sup_{x \in B} \left| \int_E g_h(S(t_0)x + y) \, d\mu_{t_0}(y) \right| \right) \\ &\leqslant 2\varepsilon \, \|f\| + \lim_{h \downarrow 0} \left(\sup_{x \in B} \int_{K_1} g_h(S(t_0)x + y) \, d\mu_{t_0}(y) \right) \\ &= 2\varepsilon \, \|f\|. \end{split}$$

For Hilbert spaces E it is known that (5.1) holds for all $f \in BUC(E)$; semigroups on BUC(E) satisfying (5.1) have been studied from an abstract point of view in [Ce] and [CG]. In our more general setting we do not know whether (5.1) holds without additional assumptions. For this reason we will impose stronger assumptions on **S** and Q.

Let $t_0 > 0$ be fixed. The pair (\mathbf{S}, Q) is said to be *null controllable at* t_0 if $S(t_0)E \subset H_{t_0}$. This condition arises in control theory in a natural way; for its interpretation and further discussion we refer to [DZ3]. If the domain D(A) of the generator A of a differentiable semigroup \mathbf{S} is contained in the RKHS H associated with Q, then (\mathbf{S}, Q) is null controllable at all t > 0; this follows from [Nv, Lemma 2.2]. Under a null controllability condition, the results of Section 4 and Theorem 5.5 are applicable. This is the content of the following proposition.

Proposition 5.7. If (\mathbf{S}, Q) is null controllable at t_0 , then:

(i) For all $t \ge t_0$ we have $\mu_t \sim \mu_{t_0}$;

(ii) For all h > 0 the operator $S_{h \to t_0+h}(t_0)$ is Hilbert-Schmidt from H_h into H_{t_0+h} ;

(iii) For all $t \ge t_0$ the operator S(t) is compact in E.

Proof: First notice that the null controllability condition implies $S(t_0)H_{t_0} \subset H_{t_0}$, so that $H_t = H_{t_0}$ for all $t \in [t_0, \infty)$, and for each h > 0, $S_{h \to t_0+h}(t_0)$ is a strict contraction.

If we regard $S(t_0)$ as an element of $\mathcal{L}(E, H_{t_0})$, then $S_{h \to t_0+h}(t_0)$ admits the factorization $S_{h \to t_0+h}(t_0) = j_{t_0 \to t_0+h} \circ S(t_0) \circ i_h$, where $i_h : H_h \subset E$ and $j_{t_0 \to t_0+h} :$ $H_{t_0} \subset H_{t_0+h}$ are the inclusion maps. By a result of Kwapień and Szymański [KS], there exists an orthonormal basis (g_n) of H_h such that $\sum_{n=1}^{\infty} ||i_h g_n||_E^2 < \infty$. But then also

$$\sum_{n=1}^{\infty} \|S_{h\to t_0+h}(t_0)g_n\|_{H_{t_0+h}}^2 \leqslant \|j_{t_0\to t_0+h} \circ S(t_0)\|_{\mathcal{L}(E,H_{t_0+h})} \sum_{n=1}^{\infty} \|i_h g_n\|_E^2 < \infty,$$

proving that $S_{h \to t_0+h}(t_0)$ is Hilbert-Schmidt.

The last assertion follows from the fact that by assumption $S(t_0)$ factors through H_{t_0} and the general fact from the theory of abstract Wiener spaces (cf. [Ku, Section 1.4]) that the inclusion map $i_{t_0} : H_{t_0} \subset E$ is compact.

Corollary 5.8. Let $t_0 > 0$ be fixed and suppose the pair (\mathbf{S}, Q) is null controllable at t_0 . Then for all bounded sets $B \subset E$ and all $f \in BUC(E)$ we have

$$\lim_{h \downarrow 0} \left(\sup_{x \in B} |P(t_0 + h)f(x) - P(t_0)f(x)| \right) = 0.$$

Proof: The null controllability assumption $S(t_0)E \subset H_{t_0}$ implies that $\mu_t \sim \mu_{t_0}$ for $t \in [t_0, \infty)$, and that for all h > 0 the operator $S_{h \to t_0+h}(t_0)$ is Hilbert-Schmidt.

Fix
$$f \in BUC(E)$$
, $x \in E$, and $t_1 > t_0$ arbitrary. Then for $h \in [0, t_1 - t_0]$

$$P(t_0 + h)f(x) - P(t_0)f(x) = \int_E f(S(t_0 + h)x + y)(g_{t_0 + h}(y) - g_{t_0}(y)) d\mu_{t_1}(y) + \int_E f(S(t_0 + h)x + y) - f(S(t_0)x + y) d\mu_{t_0}(y).$$

As $h \downarrow 0$, by Theorem 5.5 the first integral tends to 0, uniformly in x. In order to estimate the second integral, we note that $S(t_0)$ is compact by Proposition 5.7 (iii). If $B \subset E$ is a given bounded set, it then follows from the uniform continuity of f and the strong continuity of **S** that

$$\lim_{h \downarrow 0} \left(\sup_{x \in B} \left| \int_E f(S(t_0 + h)x + y) - f(S(t_0)x + y) \, d\mu_{t_0}(y) \right| \right) = 0.$$

This shows that $\lim_{h\downarrow 0} P(t_0 + h)f(x) - P(t_0)f(x) = 0$, uniformly for $x \in B$.

The following example shows that the convergence is generally not uniformly on E, even if E is one-dimensional.

Example 5.9. Let $E = \mathbb{R}$, Q = I, and $S(t) = e^{-t}$. Then

$$\int_{\mathbb{R}} \exp(-i(e^{-t}s+\tau)) \, d\mu_t(\tau) = \exp(-i(e^{-t}s))\hat{\mu}_t(1) = (1-e^{-2t})\exp(-i(e^{-t}s)).$$

Hence, for $f(s) := \cos s$ we have

$$P(t)f(s) = (1 - e^{-2t})\cos(e^{-t}s),$$

from which we deduce that $||P(t_0 + h)f - P(t_0)f|| = 2$ for all $t_0 > 0$ and h > 0.

Remark 5.10. Strong continuity in BUC(E) with E a Hilbert space was investigated in [DL], where it was shown that for a given $f \in BUC(E)$ we have $\lim_{h \downarrow 0} ||P(h)f - f|| = 0$ if and only if

$$\lim_{h \downarrow 0} \left(\sup_{x \in E} \left| f(S(h)x) - f(x) \right| \right) = 0.$$

6. The reproducing kernel Hilbert space H_{∞}

In this section we will discuss some versions of the previous results assuming that an invariant measure μ_{∞} exists.

We return to the cylindrical setting in an arbitrary real Banach space E, i.e. Assumption 2.1 is *not* adopted and E need not be separable. Instead, will make the following

Assumption 6.1. The strong limit (in E)

$$Q_{\infty}x^* := \lim_{t \to \infty} Q_t x^*$$

exists for all $x^* \in E^*$ and defines a bounded linear operator $Q_{\infty} \in \mathcal{L}(E^*, E)$.

It is clear that the operator Q_{∞} defined in this way is positive symmetric; its RKHS is denoted by H_{∞} , and the inclusion map $H_{\infty} \subset E$ is denoted by i_{∞} . The proof of Proposition 1.3 extends to show that $H_t \subset H_{\infty}$ for all t > 0.

Theorem 6.2. For all s > 0 we have $S(s)H_{\infty} \subset H_{\infty}$, and **S** restricts to a C_0 -contraction semigroup \mathbf{S}_{∞} on H_{∞} .

Proof: The invariance of H_{∞} is proved by repeating the proof of Theorem 1.4 with t replaced by ∞ ; this also gives contractivity. It remains to prove strong continuity of \mathbf{S}_{∞} on H_{∞} .

For all $h \in H_{\infty}$ and $x^* \in E^*$ we have

$$\lim_{t\downarrow 0} [S_{\infty}(t)h, Q_{\infty}x^*]_{H_{\infty}} = \lim_{t\downarrow 0} \langle S(t)h, x^* \rangle = \langle h, x^* \rangle = [h, Qx^*]_{H_{\infty}}.$$

But \mathbf{S}_{∞} being uniformly bounded on H_{∞} , the linear subspace H_{∞}^{0} of all $g \in H_{\infty}$ such that $\lim_{t\downarrow 0} [S_{\infty}(t)h, g]_{H_{\infty}} = [h, g]_{H_{\infty}}$ is closed. Therefore, $H_{\infty}^{0} = H_{\infty}$ and \mathbf{S}_{∞} is weakly continuous. By a standard result from semigroup theory [Pa, Theorem 2.1.4], this implies that \mathbf{S}_{∞} is strongly continuous.

Under the assumption that E is a Hilbert space and Q_{∞} is trace class, this result is due to Chojnowska-Michalik and Goldys [CG3, Proposition 1] (see also [CG2, Lemma 4]). Our proof is a modification of the proof of [CG2]. In fact, an analysis of this proof led us to the discovery of Theorem 1.4.

Theorem 6.3. Let $t_0 > 0$. Then $H_{t_0} = H_{\infty}$ if and only if $||S_{\infty}(t_0)||_{H_{\infty}} < 1$. In this case, $H_{t_0} = H_t = H_{\infty}$ for all $t \in (t_0, \infty)$.

Proof: We only need to prove that $H_{\infty} \subset H_{t_0}$ if and only if $\|S_{\infty}(t_0)\|_{H_{\infty}} < 1$.

We note that

$$S_{\infty}^{*}(t)Q_{\infty} = (i_{\infty}S_{\infty}(t))^{*} = (S(t)i_{\infty})^{*} = Q_{\infty}S^{*}(t);$$
(6.1)

here $i_{\infty}: H_{\infty} \to E$ is the inclusion map. First assume $||S_{\infty}(t_0)||_{H_{\infty}} < 1$. Using (6.1), for all $x^* \in E^*$ we have

$$\begin{aligned} \|Q_{t_0}x^*\|_{H_{t_0}}^2 &= \langle Q_{\infty}x^*, x^* \rangle - \langle S(t_0)Q_{\infty}S^*(t_0)x^*, x^* \rangle \\ &= \|Q_{\infty}x^*\|_{H_{\infty}}^2 - \|Q_{\infty}S^*(t_0)x^*\|_{H_{\infty}}^2 \\ &= \|Q_{\infty}x^*\|_{H_{\infty}}^2 - \|S_{\infty}^*(t_0)Q_{\infty}x^*\|_{H_{\infty}}^2 \\ &\geqslant \left(1 - \|S_{\infty}(t_0)\|_{H_{\infty}}^2\right) \|Q_{\infty}x^*\|_{H_{\infty}}^2 \end{aligned}$$

This gives the inclusion $H_{\infty} \subset H_{t_0}$.

The converse follows from an obvious modification of the proof of Theorem 1.7.

Under the assumption that E is Hilbert and Q_{∞} is trace class, this result was obtained in the second part of [CG2, Lemma 4], with a similar proof. In fact, this motivated our Theorem 1.7.

The following result gives a criterion for equality $H_{t_0} = H_{\infty}$ in terms of mapping properties of **S**.

Theorem 6.4. If $S(t_0)H_{\infty} \subset H_{t_0}$, then $H_{t_0} = H_t = H_{\infty}$ for all $t \in [t_0, \infty]$.

Proof: We always have $H_{t_0} \subset H_t \subset H_\infty$, so we only need prove the inclusion $H_\infty \subset H_{t_0}$.

First note that for all $x^* \in E^*$,

$$Q_{\infty}x^* = Q_{t_0}x^* + S(t_0)(Q_{\infty}S^*(t_0)x^*) \in H_{t_0}.$$

Next fix $h \in H_{\infty}$ arbitrary. Let $(x_n^*) \subset E^*$ be a sequence such that $\lim_{n\to\infty} Q_{\infty}x_n^* = h$ in H_{∞} . Then

$$\lim_{n \to \infty} S(t_0) Q_{\infty} S^*(t_0) x_n^* = \lim_{n \to \infty} S(t_0) S_{\infty}^*(t_0) Q_{\infty} x_n^* = S(t_0) S_{\infty}^*(t_0) h =: g$$

in H_{∞} . Note that $g \in H_{t_0}$ by the assumption on $S(t_0)$. Moreover, in H_{∞} we have

$$\lim_{n \to \infty} Q_{t_0} x_n^* = \lim_{n \to \infty} (Q_\infty x_n^* - S(t_0) Q_\infty S^*(t_0) x_n^*) = h - g$$

On the other hand, from $||Q_{t_0}x_n^*||_{H_{t_0}} \leq ||Q_{\infty}x_n^*||_{H_{\infty}}$ we see that the sequence $(Q_{t_0}x_n^*)$ is bounded in H_{t_0} . Let y be a weak limit point of $(Q_{t_0}x_n^*)$ in H_{t_0} . By the continuity of the inclusion $H_{t_0} \subset H_{\infty}$, y is also a weak limit point of $(Q_{t_0}x_n^*)$ in H_{∞} . Therefore we must have y = h - g. In particular, $h - g \in H_{t_0}$. But then $h = y + g \in H_{t_0}$.

For E Hilbert and Q_{∞} trace class, this is proved in [CG3, Proposition 3] by control theoretic methods.

The following example, taken from [Go], shows that it may happen that $H_t = H_s$ for all $t, s \in (0, \infty)$, although the inclusions $H_t \subset H_\infty$ are strict. In [Go] these facts are checked by explicit calculations; here, we derive them as consequences of our abstract results and as such the example serves as an interesting illustration of them.

Example 6.5. Let $E = l^2$ and denote by (e_n) the standard unit basis of E. Define $Q \in \mathcal{L}(E)$ by $Qe_n := e_n/n^3$. Then Q is a non-negative self-adjoint trace class operator and hence the covariance of a Gaussian measure μ on E. Define the operator A by $Ae_n := -e_n/n$. Then A is bounded on E and $S(t) := e^{tA}$ defines a uniformly continuous semigroup of self-adjoint operators on E satisfying ||S(t)|| = 1 for all $t \ge 0$.

Fix t > 0. It is easy to check that

$$Q_t = \frac{A^2}{2}(1 - S(2t)),$$
$$Q_\infty = \frac{A^2}{2}.$$

Since A^2 and S(t) commute, so do Q_t and S(t) and we see that S(t) maps $\operatorname{Im} Q_t$ into itself. We check that S(t) extends to a bounded operator on H_t . For all $h \in E$ of the form $h = \sum_{k=1}^n a_k e_k$ we have

$$\begin{split} \|S(t)Q_th\|_{H_t}^2 &= \|Q_tS(t)h\|_{H_t}^2 \\ &= [Q_tS(t)h, S(t)h]_E \\ &= [Q_th, S(2t)h]_E \\ &= \sum_{k=1}^n a_k^2 \cdot e^{-2t/k} \cdot \frac{1}{2k^2} (1 - e^{-2t/k}) \\ &\leqslant \sum_{k=1}^n a_k^2 \cdot \frac{1}{2k^2} (1 - e^{-2t/k}) \\ &= [h, Q_th]_E = \|Q_th\|_{H_t}^2. \end{split}$$

Since the set of all $Q_t h$, with h of the above form, is dense in H_t , this shows that the restriction of S(t) to $\operatorname{Im} Q_t$ extends to a contraction on H_t . Theorem 1.9 now shows that $H_t = H_s$ for all $t, s \in (0, \infty)$. On the other hand, S(t) also commutes with Q_{∞} and for t > 0 fixed we have

$$\begin{split} \|S_{\infty}(t)Q_{\infty}e_{n}\|_{H_{\infty}}^{2} &= \|Q_{\infty}S(t)e_{n}\|_{H_{\infty}}^{2} \\ &= [Q_{\infty}S(t)e_{n}, S(t)e_{n}]_{E} \\ &= e^{-2t/n}[Q_{\infty}e_{n}, e_{n}]_{E} \\ &= e^{-2t/n}\|Q_{\infty}e_{n}\|_{H_{\infty}}^{2}. \end{split}$$

Hence, $||S_{\infty}(t)||_{H_{\infty}} \ge e^{-t/n}$ for all n, so $||S_{\infty}(t)||_{H_{\infty}} = 1$. Hence by Theorem 6.3, the inclusion $H_t \subset H_{\infty}$ is strict.

Finally a simple computation shows that for all $t_0 > 0$, the restriction of $S(t_0)$ to H_{t_0} fails to be Hilbert-Schmidt. Hence, $\mu_t \perp \mu_s$ for all $t \neq s \in (0, \infty)$ by Corollary 3.6.

For the rest of this section, E is assumed to be separable and we will assume the following simultaneous strengthening of Assumptions 2.1 and 6.1:

Assumption 6.6. Assumption 6.1 holds and the cylindrical measure μ_{∞} associated with Q_{∞} is countably additive.

In other words, we assume that the operator Q_{∞} is the covariance of a centered Gaussian measure μ_{∞} on the Borel σ -algebra of E.

Remark 6.7. The following conditions are sufficient for Assumption 6.6 to hold:

- (i) E is a Hilbert space and $\sup_{t>0} \operatorname{Trace} Q_t < \infty$ [DZ3, Chapter 11];
- (ii) E is a Hilbert space, Q is trace class, and **S** is uniformly exponentially stable;
- (iii) The cylindrical measure associated with Q is countably additive, **S** is uniformly exponentially stable, $S(s)H \subset H$ for all $s \ge 0$, and

$$\int_0^\infty \|S(s)\|_{\mathcal{L}(H)}^2 \, ds < \infty$$

(iv) Assumption 2.1 holds, **S** is uniformly exponentially stable, and the pair (**S**, Q) is null controllable at some $t_0 > 0$.

We will investigate the question under what conditions we have equivalence $\mu_{t_0} \sim \mu_{\infty}$ holds for a given $t_0 \in (0, \infty)$.

Theorem 6.8. For a fixed $t_0 > 0$, the measures μ_{t_0} and μ_{∞} are equivalent if and only if the following two conditions are satisfied:

- (i) $||S_{\infty}(t_0)||_{H_{\infty}} < 1;$
- (ii) The operator $S_{\infty}(t_0)S_{\infty}^*(t_0)$ is Hilbert-Schmidt on H_{∞} .

For Hilbert spaces E, this was proved in [CG3, Theorem 2]. By the semigroup property, this result implies:

Corollary 6.9. If $\mu_{t_0} \sim \mu_{\infty}$ for some $t_0 > 0$, then $\mu_t \sim \mu_{\infty}$ for all $t \in [t_0, \infty]$.

It is possible to give an explicit expression for the Radon-Nikodym density $d\mu_{t_0}/d\mu_{\infty}$. If $\mu_{t_0} \sim \mu_{\infty}$ for some $t_0 > 0$, then $S_{\infty}(t_0)S_{\infty}^*(t_0)$ is Hilbert-Schmidt on H_{∞} . If we assume that $S_{\infty}(t_0)$ itself is Hilbert-Schmidt we can prove more:

Theorem 6.10. Suppose we have $\mu_{t_0} \sim \mu_{\infty}$ for some $t_0 > 0$. If $S_{\infty}(t_0)$ is Hilbert-Schmidt on H_{∞} , then the Radon-Nikodym derivative $g_{t_0} := d\mu_{t_0}/d\mu_{\infty}$ is μ_{∞} -a.e. given by

$$g_{t_0}(x) = \left(\det\sqrt{I - S_{\infty}(t_0)S_{\infty}^*(t_0)}\right)^{-1} \times \\ \times \exp\left(-\frac{1}{2}\|\overline{\left(\sqrt{I - S_{\infty}(t_0)S_{\infty}^*(t_0)}\right)^{-1}(S_{\infty}(t_0)S_{\infty}^*(t_0))^{\frac{1}{2}}x}\|^2\right).$$

Concerning continuous dependence, we have:

Theorem 6.11. Under the above assumptions, the Radon-Nikodym derivative $g_t := d\mu_t/d\mu_\infty$ exists for all $t \ge t_0$ and belongs to $L^2(E, \mu_\infty)$. The function $t \mapsto g_t$ is continuous from $[t_0, \infty)$ into $L^2(E, \mu_\infty)$.

Analogously to the situation encountered in Section 5, the assumptions of the theorem are automatically satisfied under the null controllability assumption $S(t_0)E \subset H_{t_0}$. The proofs of Theorems 6.10 and 6.11 proceed as in Sections 4 and 5, respectively. The main ingredient of Theorem 6.10 is the following version of Theorem 4.1:

Theorem 6.12. The semigroup **P** extends to a C_0 -semigroup on $L^2(E, \mu_{\infty})$ and for all $t \ge 0$ we have

$$P(t) = \Gamma(S_{\infty}^{*}(t)).$$

For Hilbert spaces E, Theorems 6.10 and 6.12 are due to Chojnowska-Michalik and Goldys [CG2], [CG3]. Their version of Theorem 6.10 is based on a very general formula for Radon-Nikodym derivatives of Gaussian measures on Hilbert spaces due to Fuhrman [Fu], who obtained the Hilbert space case of Theorem 6.10 under the null controllability assumption $S(t)E \subset H_t$ for all t > 0.

7. Extension to Gaussian Mehler semigroups

In [BRS], Bogachev, Röckner, and Schmuland introduced the concept of a generalized Mehler semigroup. Under Assumption 2.1, the Ornstein-Uhlenbeck semigroups **P** belong to this class. In this final section we will discuss briefly some extensions of our results to this more general framwork.

Let E be a separable real Banach space, let **S** be a C_0 -semigroup on E, and let $\{\mu_t\}_{t\geq 0}$ be a one-parameter family of probability measures defined on the Borel σ -algebra of E. The pair (**S**, $\{\mu_t\}_{t\geq 0}$) is called a *Mehler semigroup* on E if

$$\mu_{t+s} = (T(s)\mu_t) * \mu_s, \qquad t, s \ge 0, \tag{7.1}$$

where $T(s)\mu_t$ denotes the image measure of μ_t under T(s). This terminology is explained by the observation [BRS, Proposition 2.2] that $(\mathbf{S}, \{\mu_t\}_{t \ge 0})$ is a Mehler semigroup if and only if

$$P(t)f(x) := \int_E f(S(t)x - y) \, d\mu_t(y), \qquad t \ge 0, \, x \in E,$$

defines a semigroup on the space $B_b(E)$ of bounded Borel functions on E. More generally, a pair $(\mathbf{S}, \{\mu_t\}_{t\geq 0})$, where \mathbf{S} is a C_0 -semigroup on E and $\{\mu_t\}_{t\geq 0}$ is a oneparameter family of cylindrical probability measures on the ring of cylindrical sets in E, is called a *cylindrical Mehler semigroup* on E if (7.1) holds. If $Q \in \mathcal{L}(E^*, E)$ is a positive and symmetric operator and **S** is a C_0 -semigroup on E, then the pair (**S**, $\{\mu_t\}_{t \ge 0}$), where μ_t is the unique cylindrical measure whose Fourier transform is given by

$$\hat{\mu}_t(x^*) = \exp\left(-\frac{1}{2}\langle Q_t x^*, x^*\rangle\right), \qquad x^* \in E^*$$
(7.2)

is easily seen to be a cylindrical Mehler semigroup; it is a Mehler semigroup if Assumption 2.1 holds.

Motivated by this example, we say that $(\mathbf{S}, \{\mu_t\}_{t \ge 0})$ is *Gaussian* if for each t > 0there exists a positive symmetric operator $Q_t \in \mathcal{L}(E^*, E)$, the *covariance* of μ_t , such that the Fourier transform of μ_t is given by (7.2). In this situation we denote by H_t the RKHS associated with the covariance operator Q_t of μ_t . By considering the Fourier transform of (7.1) we have the identity [BRS, Proposition 2.2]

$$Q_{t+s} = Q_s + S(s)Q_t S^*(s), \qquad t, s \ge 0.$$
 (7.3)

In particular, $\langle Q_{t+s}x^*, x^* \rangle = \langle Q_sx^*, x^* \rangle + \langle Q_tS^*(s)x^*, S^*(s)x^* \rangle$ for all $t, s \ge 0$ and $x^* \in E^*$. By positivity, this shows that the functions $t \mapsto \langle Q_tx^*, x^* \rangle$ are increasing. Hence,

$$H_{t_0} \subset H_{t_1}$$
 whenever $0 < t_0 < t_1 < \infty$. (7.4)

Inspection of the proofs shows that (7.3) and (7.4) are all that is needed for most of the results in this paper. These therefore extend to Gaussian (cylindrical) Mehler semigroups without change.

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8. References

- [BRS] V.I. BOGACHEV, M. RÖCKNER, AND B. SCHMULAND, Generalized Mehler semigroups and applications, Prob. Th. Relat. Fields 105 (1996), 193-225.
 - [Ce] S. CERRAI, A Hille-Yosida theorem for weakly continuous semigroups, Semigroup Forum 49 (1994), 349-367.
- [CG] S. CERRAI AND F. GOZZI, Strong solutions of Cauchy problems associated to weakly continuous semigroups, Differential Integral Eq. 8 (1995), 465-486.
- [CG1] A. CHOJNOWSKA-MICHALIK AND B. GOLDYS, Existence, uniqueness and invariant measures for stochastic semiliniear equations on Hilbert spaces, Prob. Th. Relat. Fields 102 (1995), 331-356.

- [CG2] A. CHOJNOWSKA-MICHALIK AND B. GOLDYS, Nonsymmetric Ornstein-Uhlenbeck semigroup as second quantized operator, J. Math. Kyoto Univ. 36 (1996), 481-498.
- [CG3] A. CHOJNOWSKA-MICHALIK AND B. GOLDYS, On regularity properties of nonsymmetric Ornstein-Uhlenbeck semigroups in L^p spaces, Stoch. Stoch. Rep. 59 (1996), 183-209.
- [DL] G. DA PRATO AND A. LUNARDI, On the Ornstein-Uhlenbeck operator in spaces of continuous functions, J. Func. Anal. 131 (1995), 94-114.
- [DZ1] G. DA PRATO AND J. ZABCZYK, Smoothing properties of transition semigroups in Hilbert spaces, Stochastics 35 (1991), 63-77.
- [DZ2] G. DA PRATO AND J. ZABCZYK, Regular densities of invariant measures in Hilbert Spaces, J. Func. Anal. 130 (1995), 427-449.
- [DZ3] G. DA PRATO AND J. ZABCZYK, Stochastic Equations in Infinite Dimensions, Encyclopedia of Mathematics and its Applications, Cambridge University Press (1992).
- [DU] J. DIESTEL AND J.J. UHL, Vector Measures, Math. Surveys nr. 15, Amer. Math. Soc., Providence, R.I. (1977).
- [F-LP] D. FEYEL AND A. DE LA PRADELLE, Opérateurs linéaires gaussiens, Potential Anal. 3 (1994), 89-105.
 - [Fu] M. FUHRMAN, Densities of Gaussian measures and regularity of nonsymmetric Ornstein-Uhlenbeck semigroups in Hilbert space, preprint IMPAN no. 528 (1994).
- [GGK] I. GOHBERG, S. GOLDBERG, AND M.A. KAASHOEK, Classes of Linear Operators, Vol. I, Operator Theory: Adv. and Appl. 49, Birkhäuser (1990).
 - [Go] B. GOLDYS, On lilinear forms related to Ornstein-Uhlenbeck semigroup on Hilbert space, preprint.
 - [Gr] L. GROSS, Potential theory on Hilbert space, J. Func. Anal. 1 (1967), 123-181.
 - [Ku] H.H. KUO, Gaussian Measures on Banach Spaces, Springer Lect. Notes Math. 463, Springer-Verlag (1975).
 - [KS] S. KWAPIEŃ AND B. SZYMAŃSKI, Some remarks on Gaussian measures in Banach spaces, Prob. Math. Statistics 1 (1980), 59-65.
 - [MS] A. MILLET AND W. SMOLÉNSKI, On the continuity of Ornstein-Uhlenbeck processes in infinite dimensions, Prob. Th. Relat. Fields 92 (1992), 529-547.
 - [Nv] J.M.A.M. VAN NEERVEN, Sandwiching C_0 -semigroups, submitted for publication.
 - [NZ] J.M.A.M. VAN NEERVEN AND J. ZABCZYK, Norm discontinuity of Ornstein-Uhlenbeck semigroups, to appear in Semigroup Forum.
 - [Ne] E. NELSON, The free Markoff field, J. Funct. Anal. 1 (1973), 211-227.
 - [Pa] A. PAZY, Semigroups of Linear Operators and Applications to Partial Differential Equations, Springer-Verlag (1983).

- [Sh] D. SHALE, Linear symmetries of the free boson field, Trans. Amer. Math. Soc. 103 (1962), 149-167.
- [Si] B. SIMON, The $P(\Psi)_2$ -Euclidean (Quantum) Field Theory, Princeton University Press (1974).
- [Ta] V.I. TARIELADZE, On equivalence of Gaussian measures on Banach spaces, Bull. Acad. Sci. Georgian SSR 73 (1974), 529-532.
- [VT] N.N. VAKHANIA AND V.I. TARIELADZE, On singularity and equivalence of Gaussian measures, in: M.M. Rao (Ed.), Real and Stochastic Analysis: Recent Advances, CRC Press, Boca Raton-New York (1997), 367-389.
- [VTC] N.N. VAKHANIA, V.I. TARIELADZE, AND S.A. CHOBANYAN, Probability Distributions on Banach Spaces, D. Reidel Publishing Company, Dordrecht-Boston-Lancaster-Tokyo (1987).

Note added in proof – After this paper had been accepted for publication, the author realized that without any compactness assumption, (5.1) always holds if $f \in BUC(E)$. In fact, it turns out that one always has $\lim_{t\downarrow 0} \mu_t = \delta_0$ weakly; this is a consequence of Anderson's inequality and easily implies the assertion just made. As a consequence, in Corollary 5.8 the null controllability assumption can be omitted, and the characterization of strong continuity in BUC(E) mentioned in Remark 5.10 extends to Banach spaces E. The details will be presented elsewhere.